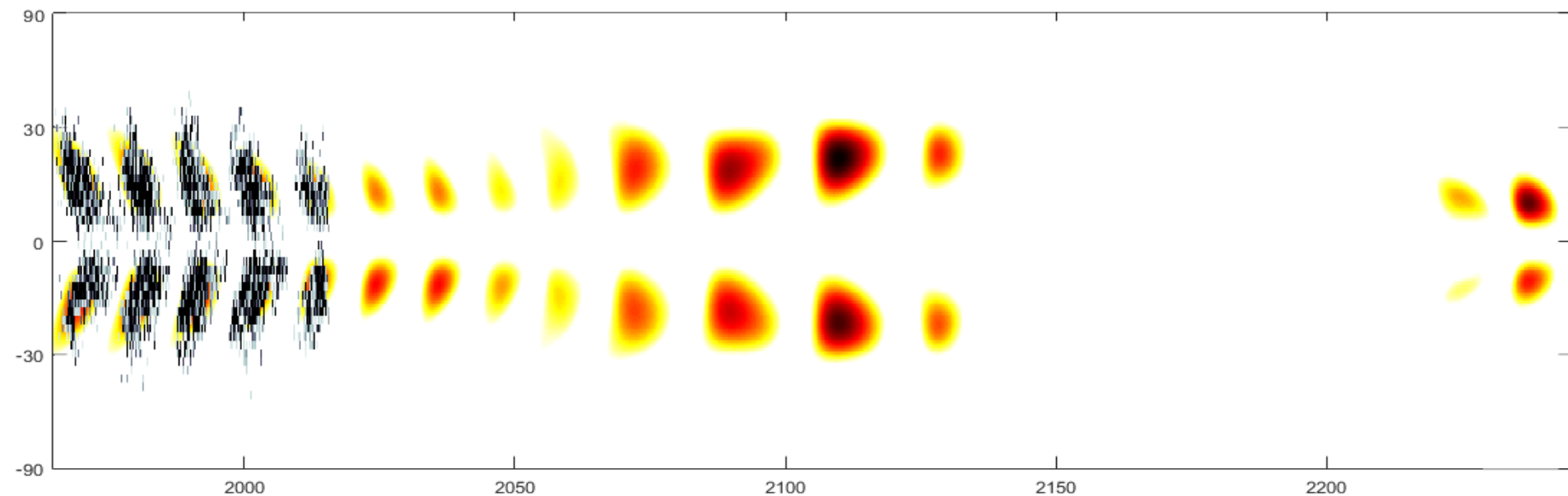
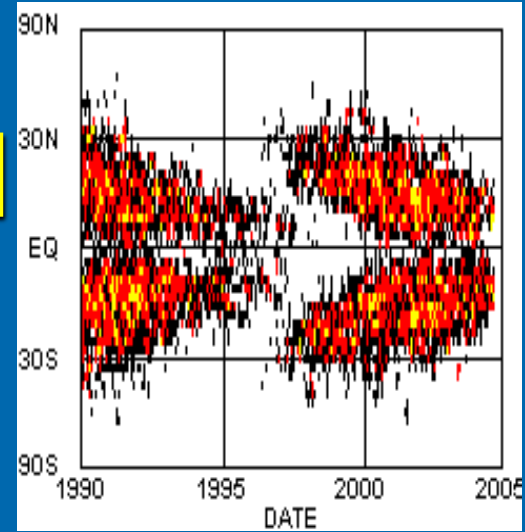
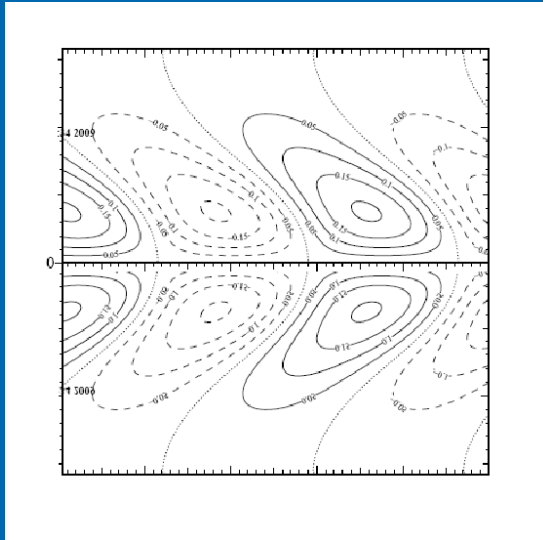


INTRODUCTION TO DYNAMO

THEORY

N. KLEEORIN

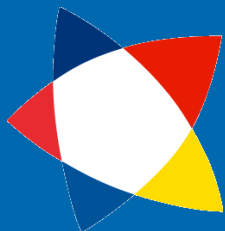


I. ROGACHEVSKI



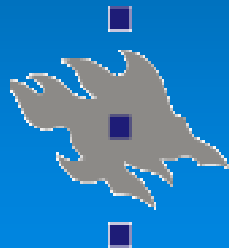
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Induction equation – from Maxwell equation

$$\nabla \cdot \mathbf{D} = 4\pi\rho_e$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{H} = \frac{4\pi}{c} \mathbf{j} + \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t}$$

$$\frac{\partial \rho_e}{\partial t} + \nabla \cdot \mathbf{j} = 0$$

Induction equation – from Maxwell equation

$$\nabla \cdot \mathbf{D} = 4\pi\rho_e$$

$$\mathbf{D} = \varepsilon\mathbf{E};$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\mathbf{B} = \mu\mathbf{H}$$

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

$$\mathbf{E}' = \mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}$$

$$\nabla \times \mathbf{H} = \frac{4\pi}{c} \mathbf{j} + \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t}$$

$$\mathbf{j} = \sigma \mathbf{E}' = \sigma \left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} \right)$$

$$\frac{\partial \rho_e}{\partial t} + \nabla \cdot \mathbf{j} = 0$$

Induction equation – from Maxwell equation

➤ Induction equation

$$\nabla \cdot \mathbf{D} = 4\pi\rho_e = 0$$

$$\mathbf{D} = \mathbf{E};$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\mathbf{B} = \mathbf{H}$$

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

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$$\nabla \times \mathbf{H} = \frac{4\pi}{c} \mathbf{j} + \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t}$$

$$\Leftarrow L_{sist} \ll cT_{sist}$$

$$\frac{\partial \rho_e}{\partial t} + \nabla \cdot \mathbf{j} = 0$$

Induction equation – from Maxwell equation

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} \quad \mathbf{E}' = \mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}$$

$$\nabla \times \mathbf{B} \approx \frac{4\pi}{c} \mathbf{j} \quad \mathbf{j} = \sigma \mathbf{E}' = \sigma \left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} \right)$$

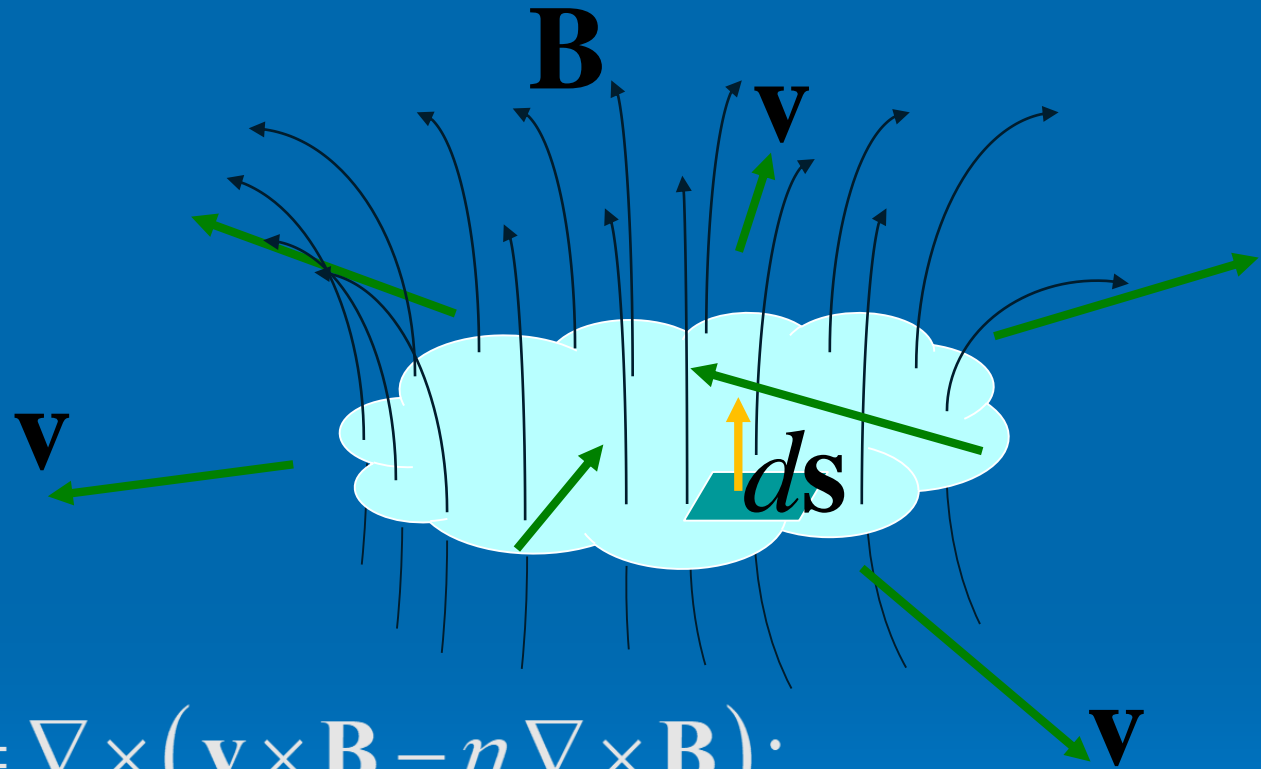
$$\mathbf{E}' = \frac{\mathbf{j}}{\sigma} \approx \frac{c \nabla \times \mathbf{B}}{4\pi\sigma} \Rightarrow \mathbf{E} = \mathbf{E}' - \frac{\mathbf{v}}{c} \times \mathbf{B} \Rightarrow \mathbf{E} \approx \frac{c \nabla \times \mathbf{B}}{4\pi\sigma} - \frac{\mathbf{v}}{c} \times \mathbf{B}$$

$$\nabla \times \left(\frac{c \nabla \times \mathbf{B}}{4\pi\sigma} - \frac{\mathbf{v}}{c} \times \mathbf{B} \right) = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B} - \eta \nabla \times \mathbf{B}); \quad \eta = \frac{c^2}{4\pi\sigma}$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \Delta \mathbf{B}$$

Induction equation – Properties

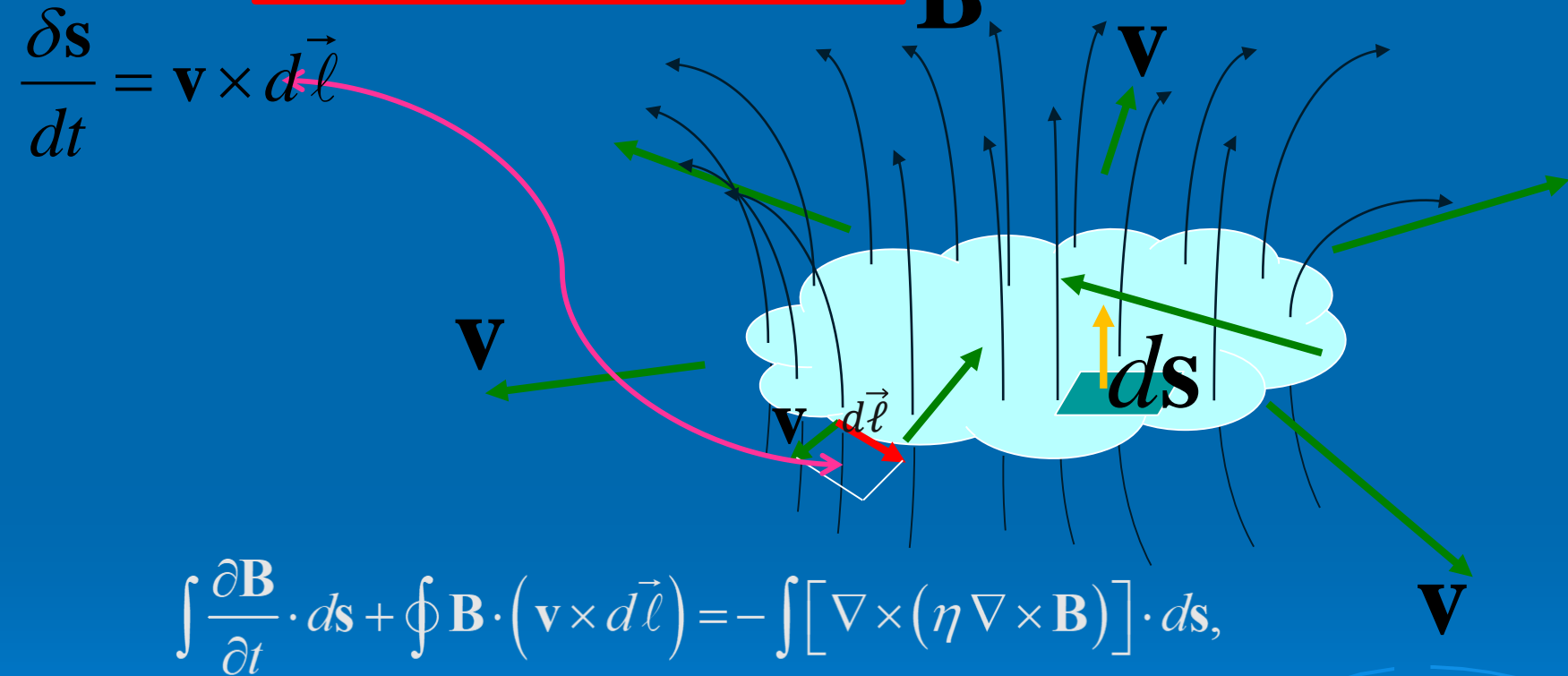


$$\int \cdot ds \left| \frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B} - \eta \nabla \times \mathbf{B}); \right.$$

$$\int \frac{\partial \mathbf{B}}{\partial t} \cdot ds = \oint (\mathbf{v} \times \mathbf{B}) \cdot d\vec{\ell} - \int \nabla \times (\eta \nabla \times \mathbf{B}) \cdot ds;$$

Induction equation – Properties

If $\eta = 0$ dynamo *impossible*



$$\int \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{s} + \oint \mathbf{B} \cdot (\mathbf{v} \times d\vec{\ell}) = - \int [\nabla \times (\eta \nabla \times \mathbf{B})] \cdot d\mathbf{s},$$

⇓

$$\int \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{s} + \oint \mathbf{B} \cdot \frac{\delta \mathbf{s}}{dt} \equiv \frac{d}{dt} \int \mathbf{B} \cdot d\mathbf{s} = \frac{d\Phi_{\mathbf{B}}}{dt} = - \int [\nabla \times (\eta \nabla \times \mathbf{B})] \cdot d\mathbf{s},$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B} - \eta \nabla \times \mathbf{B}); \quad \text{Alpha-Omega Dynamo (Mean-Field Approach)}$$

➤ Induction equation for **mean magnetic field**:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{U} \times \mathbf{B} + \boldsymbol{\varepsilon} - \eta \nabla \times \mathbf{B})$$

➤ **Electromotive force**: Shteenbeck, Krause and Raedler (1966)

$$\mathbf{v} = \mathbf{U} + \mathbf{u}; \quad \mathbf{B} \rightarrow \mathbf{B} + \mathbf{b}$$

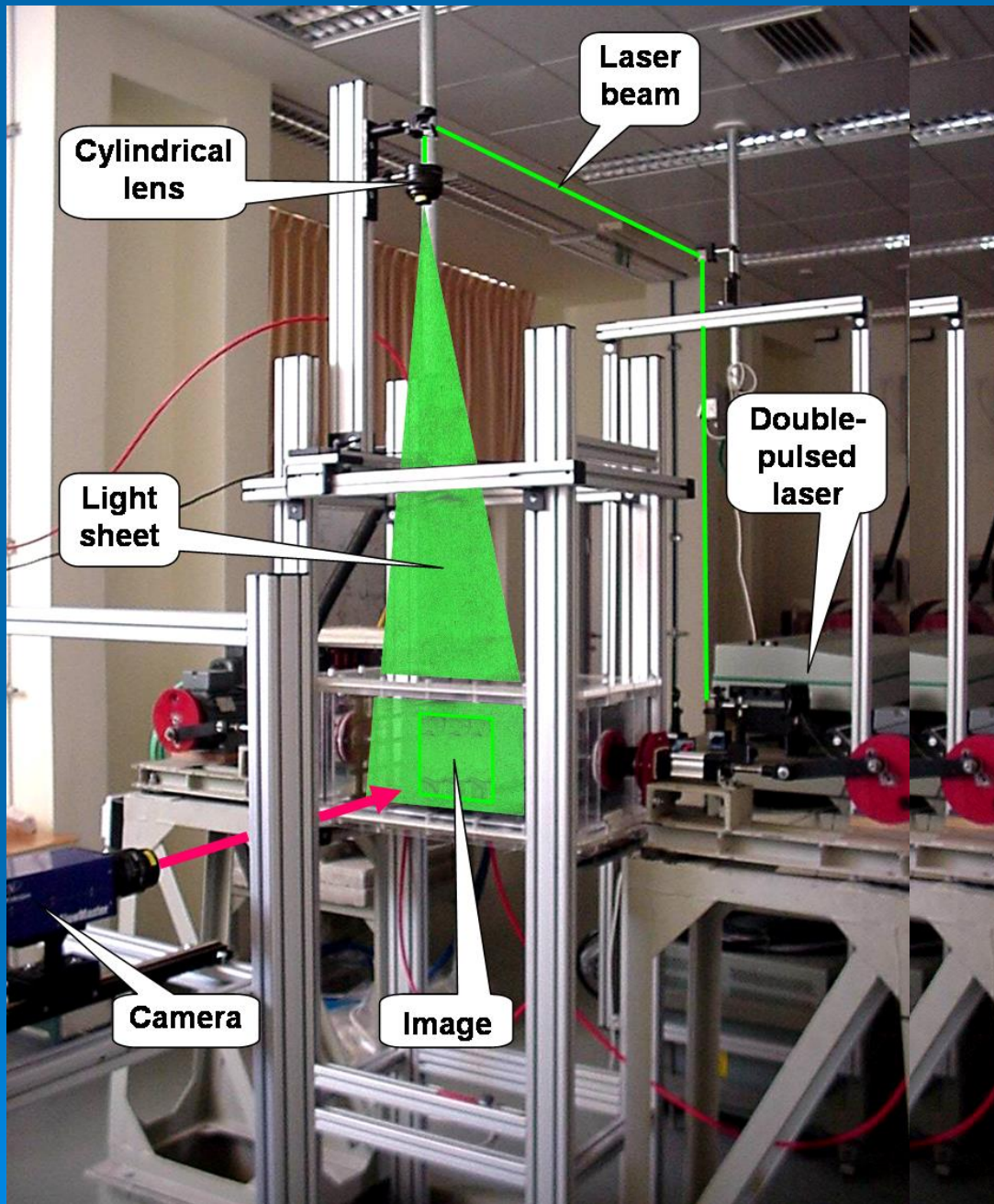
$$\boldsymbol{\varepsilon} \equiv \langle \mathbf{u} \times \mathbf{b} \rangle = \alpha \mathbf{B} - \eta_T \nabla \times \mathbf{B} + \dots$$

$$\alpha = -\frac{\tau}{3} \langle \mathbf{u} \cdot \text{rot } \mathbf{u} \rangle; \quad \eta_T = \frac{\tau \langle \mathbf{u}^2 \rangle}{3}$$

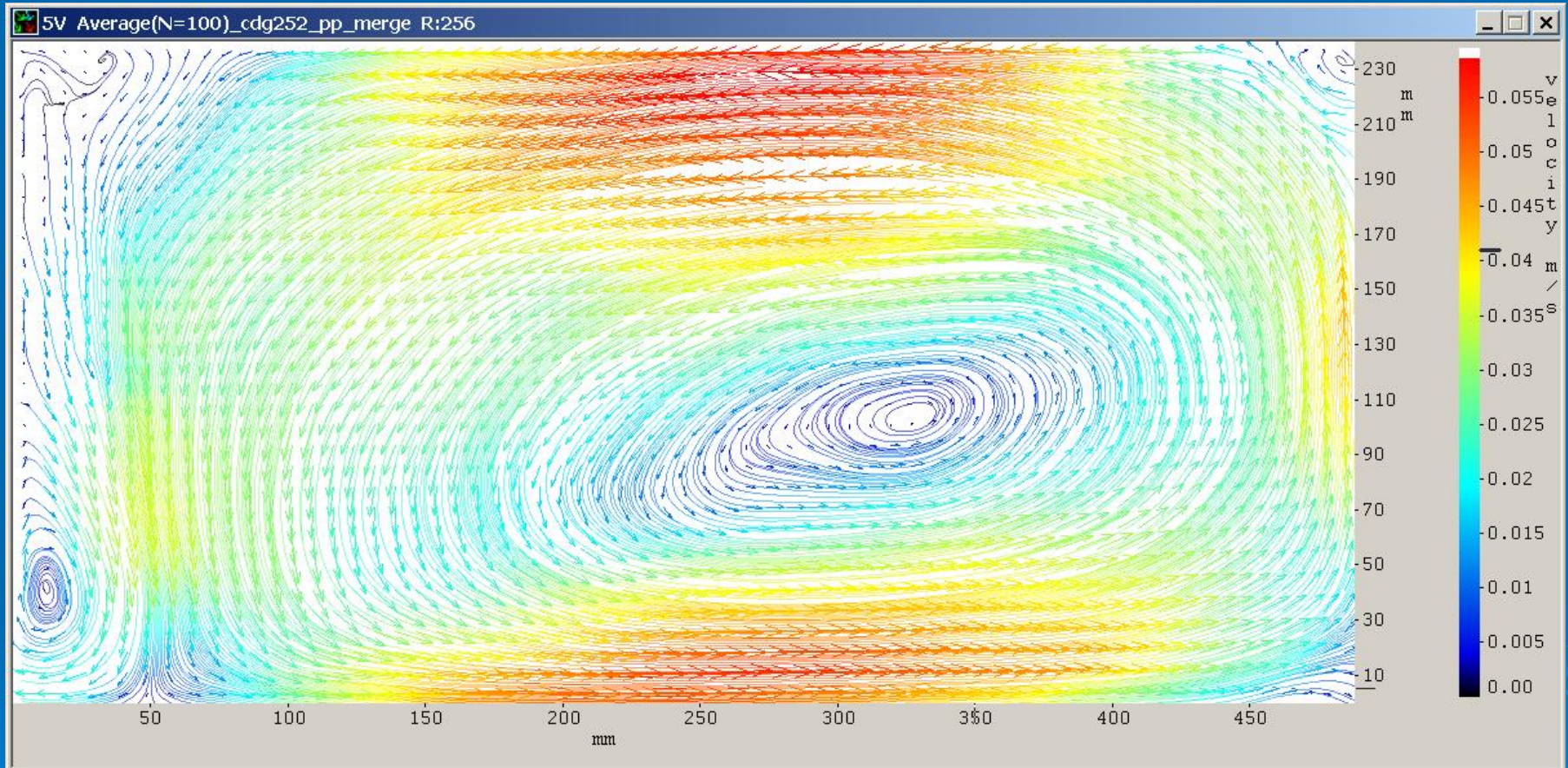


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Experimental set – up



Laboratory Turbulent Convection



Before averaging

Mean-Field Approach

- Induction equation for **mean magnetic field**:

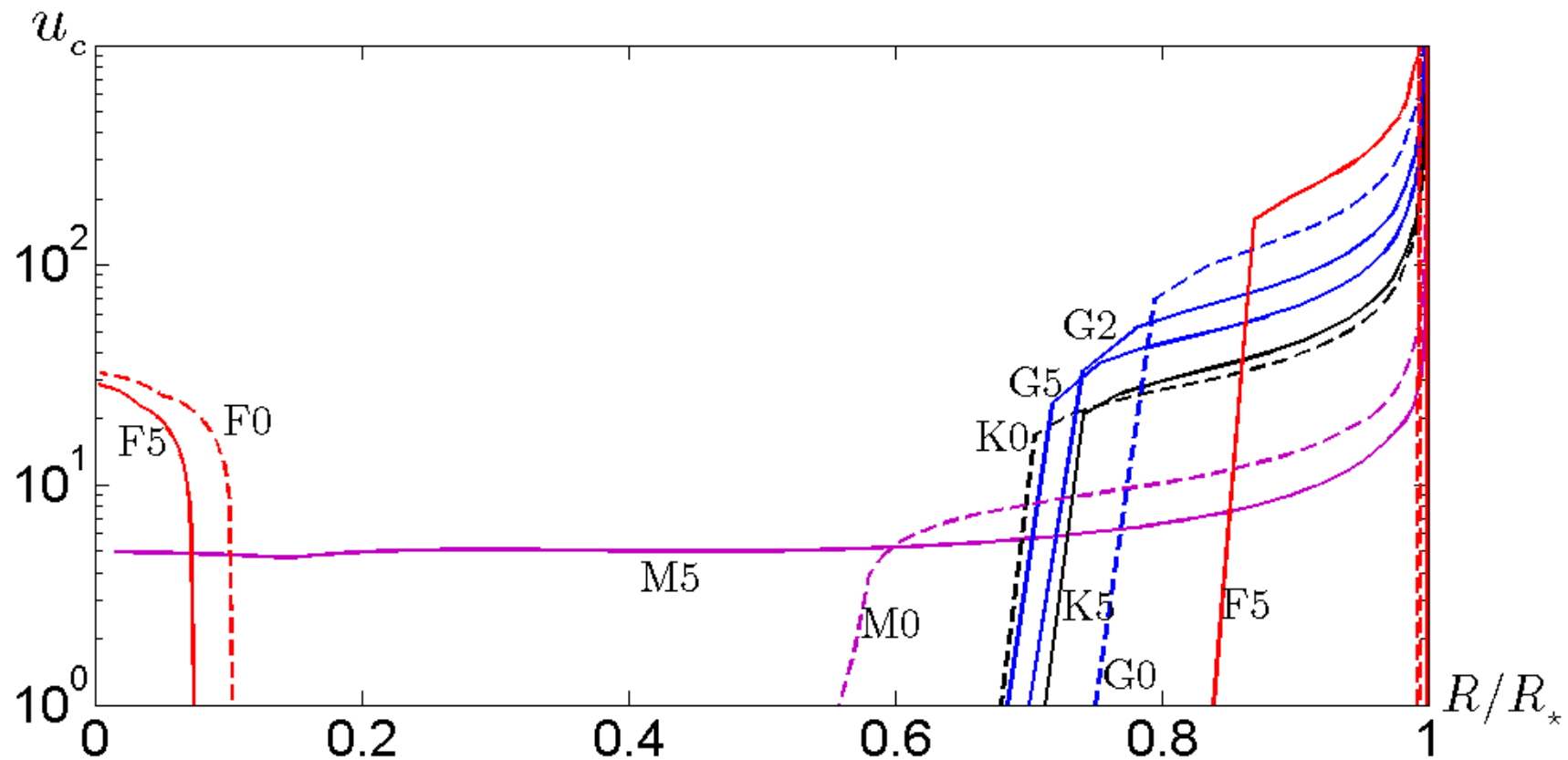
$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{U} \times \mathbf{B} + \boldsymbol{\varepsilon} - \eta \nabla \times \mathbf{B})$$

- **Electromotive force**:

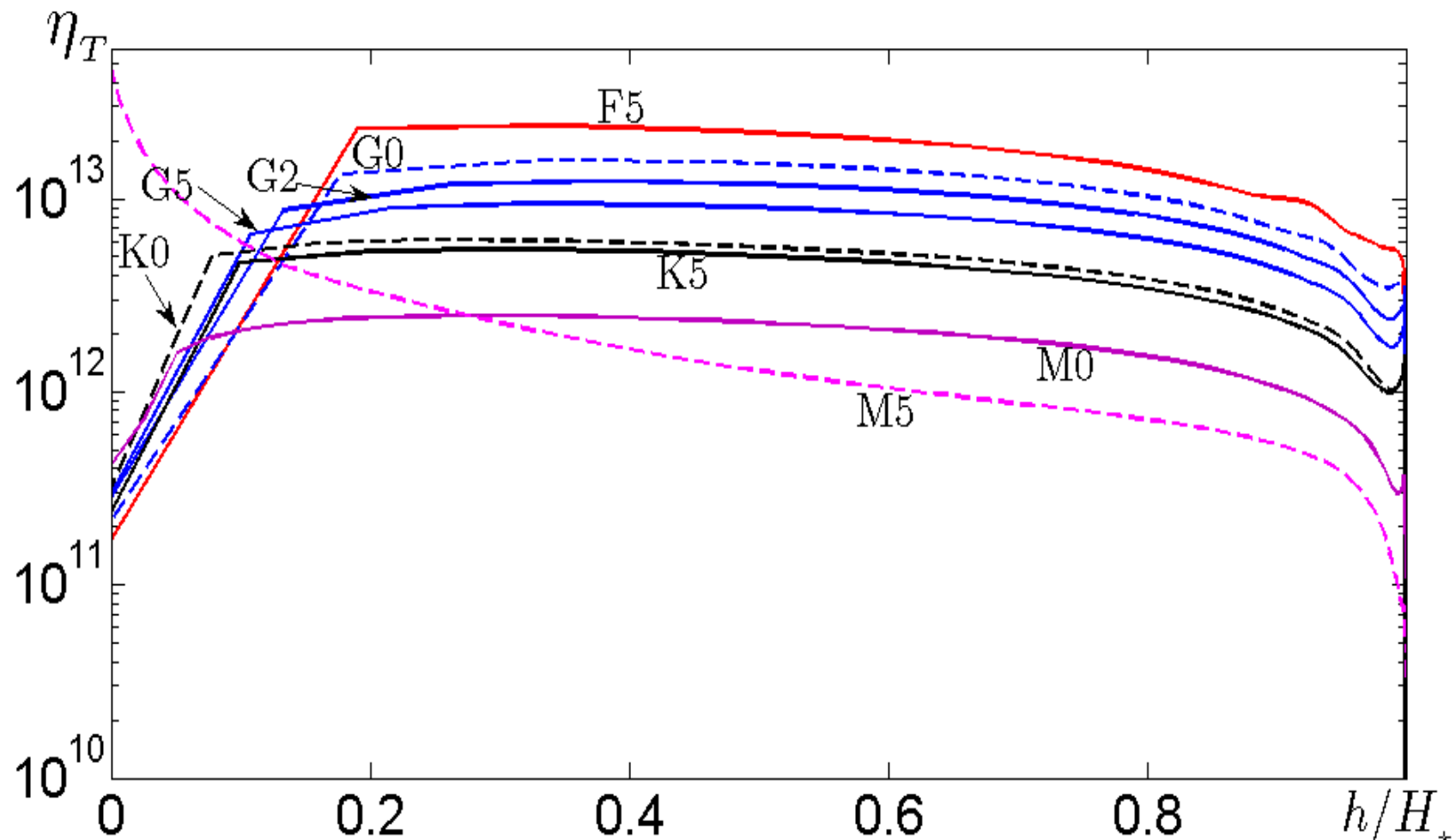
$$\boldsymbol{\varepsilon} \equiv \langle \mathbf{u} \times \mathbf{b} \rangle = \alpha \mathbf{B} - \eta_T \nabla \times \mathbf{B} + \dots$$

$$\alpha = -\frac{\tau}{3} \langle \mathbf{u} \cdot \text{rot } \mathbf{u} \rangle; \quad \eta_T = \frac{\tau}{3} \ell_0 u_c$$

Convective velocity in a star of main sequence

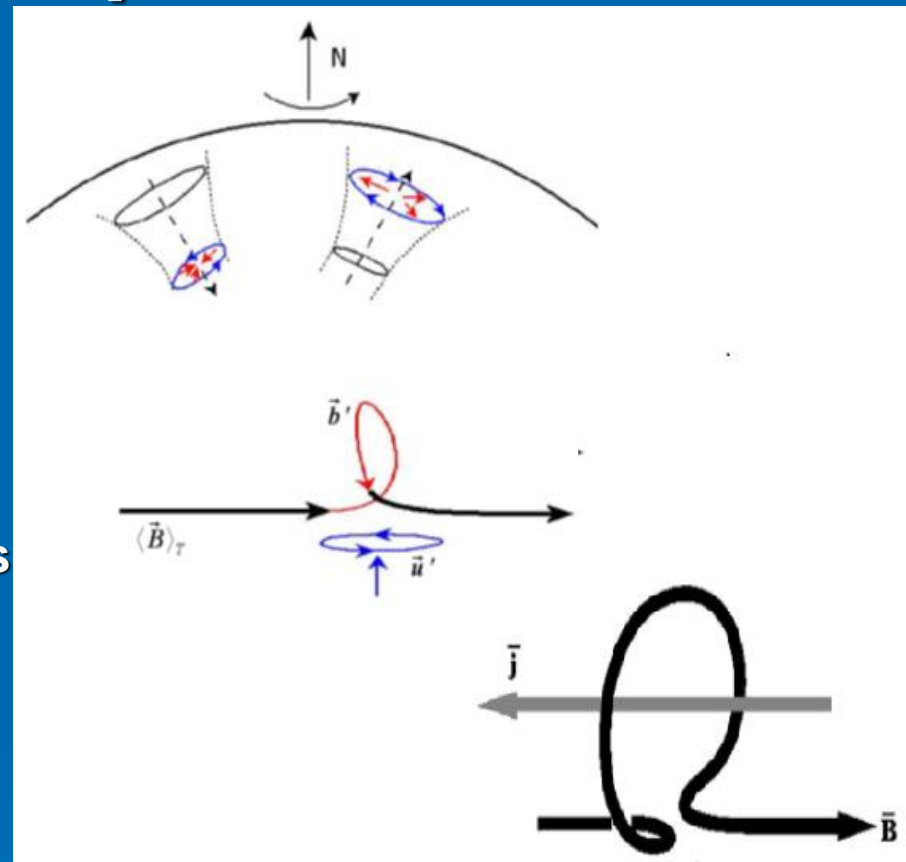


Turbulent diffusivity in a star of main sequence



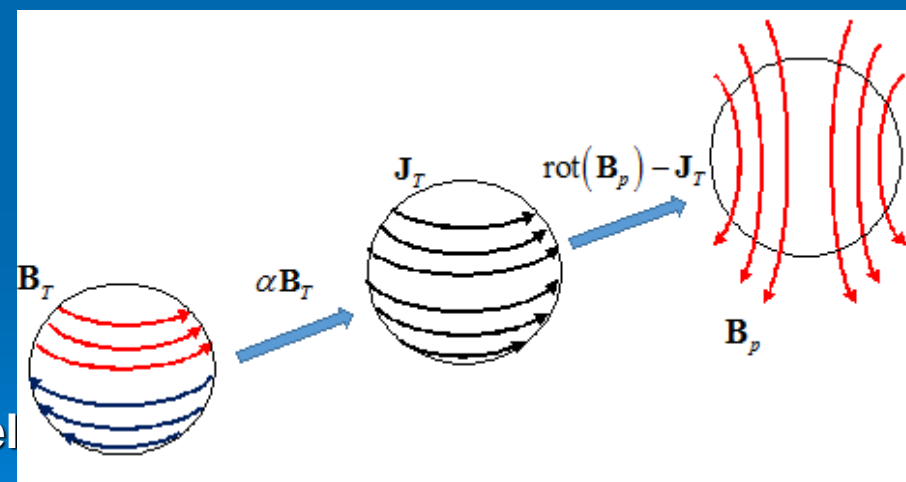
Physics of the alpha-effect

- The α -effect is related to the hydrodynamic helicity in an inhomogeneous turbulence.
- The deformations of the magnetic field lines are caused by upward and downward rotating turbulent eddies.
- The inhomogeneity of turbulence breaks a symmetry between the upward and downward eddies.
- Therefore, the total effect of the upward and downward eddies on the mean magnetic field does not vanish and it creates the mean electric current parallel to the original mean magnetic field.



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Alpha-effect in Dynamo (Mean-Field Approach)

➤ Induction equation for **mean magnetic field**:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times \left(\mathbf{U} \times \mathbf{B} + \alpha \mathbf{B} - (\eta_T + \eta) \nabla \times \mathbf{B} \right)$$

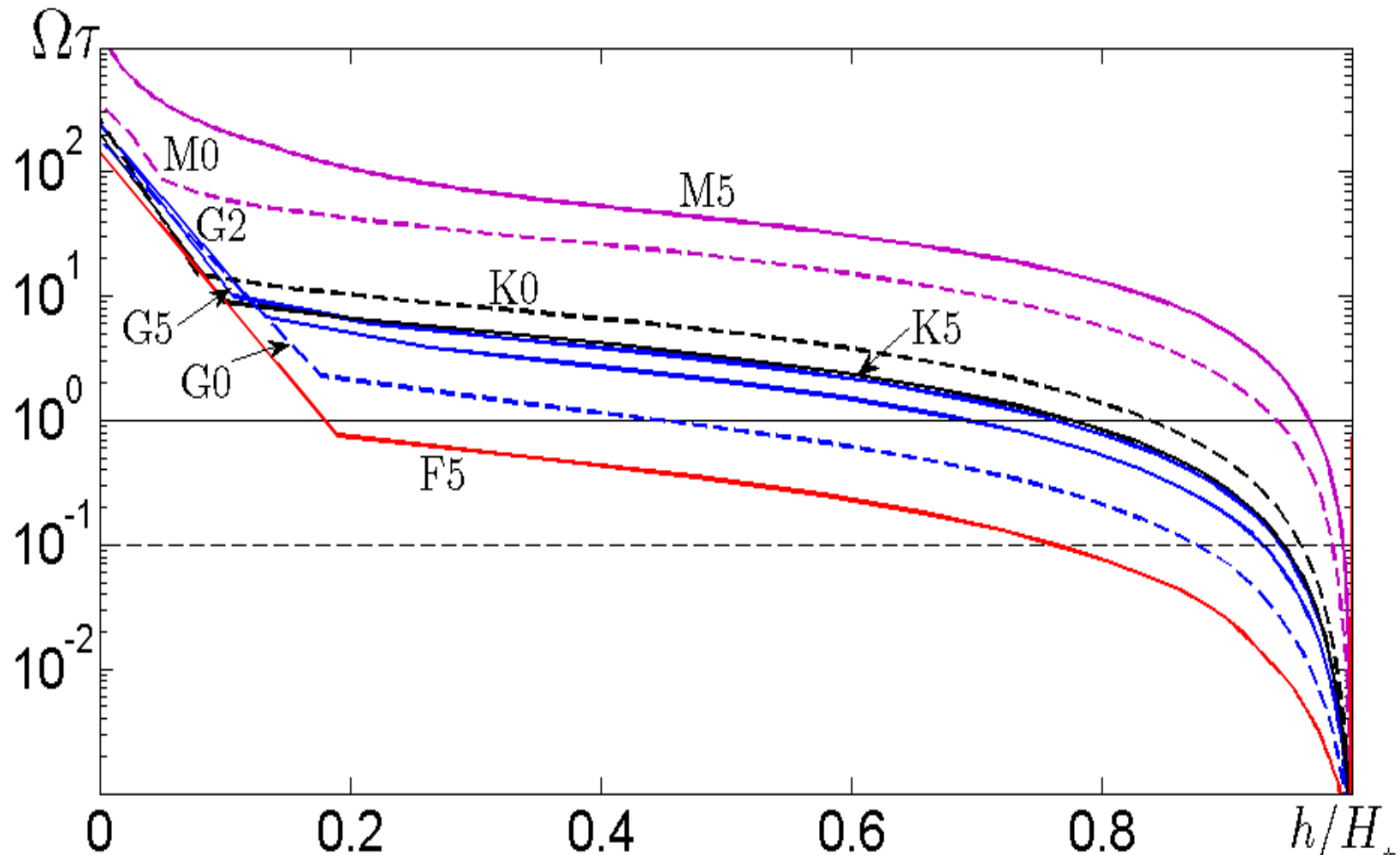
➤ **Electromotive force**:

$$\boldsymbol{\varepsilon} \equiv \langle \mathbf{u} \times \mathbf{b} \rangle = \alpha \mathbf{B} - \eta_T \nabla \times \mathbf{B} + \dots \quad \nabla \times \mathbf{B} \approx \frac{4\pi}{c} \mathbf{J}$$

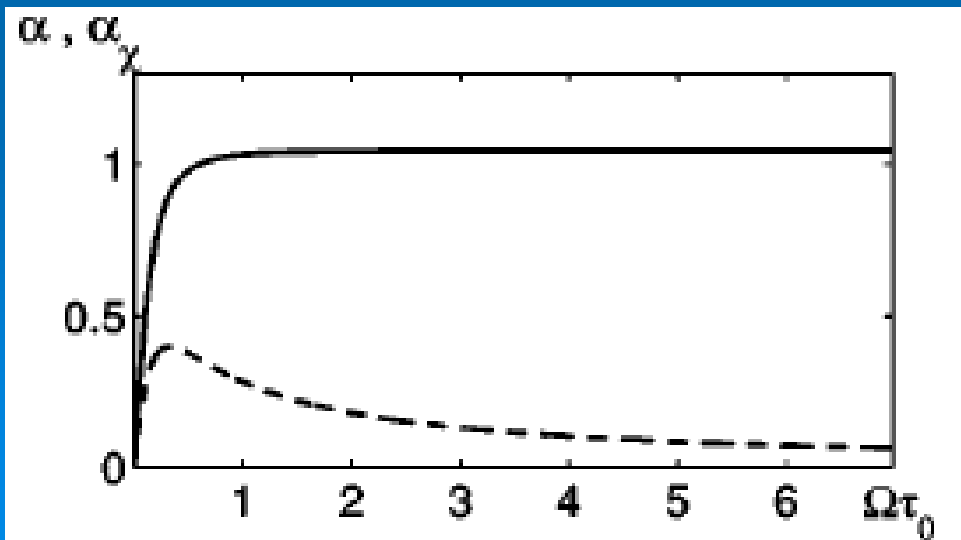
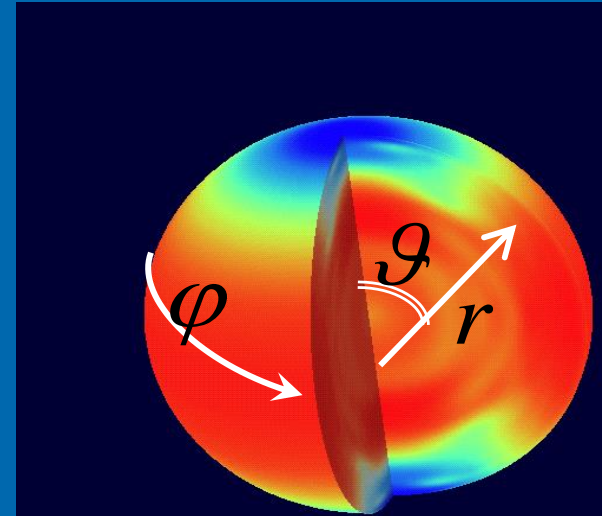
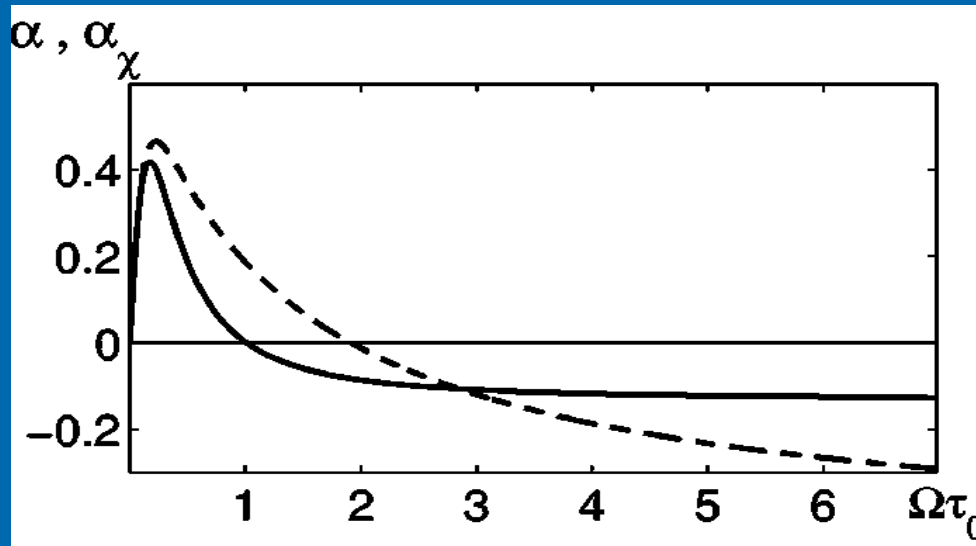
$$\mathbf{J} = \sigma_{eff} \left[\mathbf{E} + \frac{1}{c} (\mathbf{U} \times \mathbf{B} + \alpha \mathbf{B}) \right]; \quad \sigma_{eff} = \frac{c^2}{4\pi (\eta_T + \eta)}$$

$$\eta_T = \frac{\ell_0 u_c}{3}; \quad \alpha = -\frac{\tau}{3} \langle \mathbf{u} \cdot \mathbf{rot} \mathbf{u} \rangle \approx \frac{\Omega \ell_0^2}{3L_\rho} \cos(\vartheta); \quad L_\rho^{-1} = -\frac{1}{\rho} \frac{\partial \rho}{\partial r}, \quad \Omega \tau \ll 1$$

Rossby-parameter $\Omega_* \tau_*$ for main sequence stars



Radial profiles of the alpha-effect



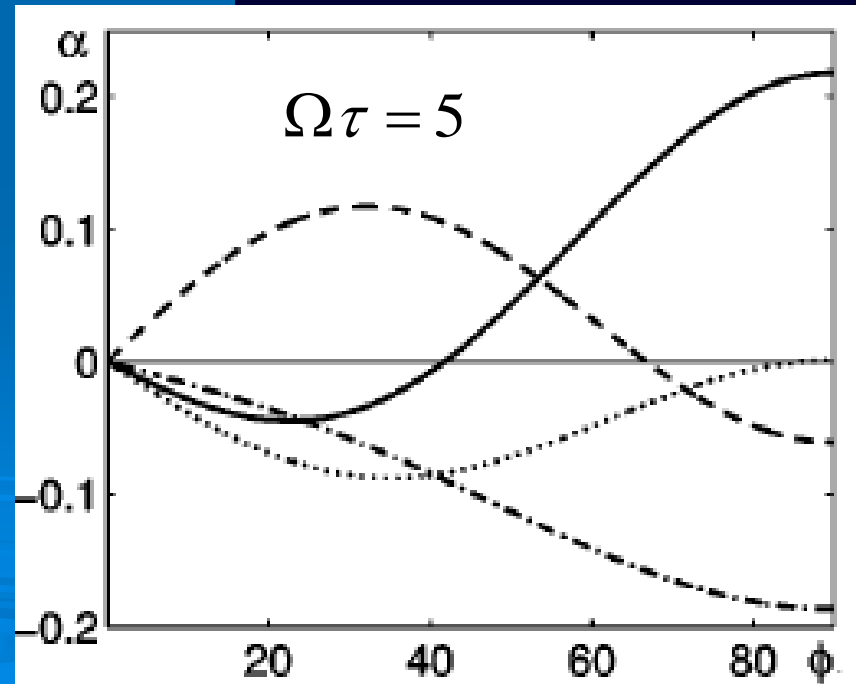
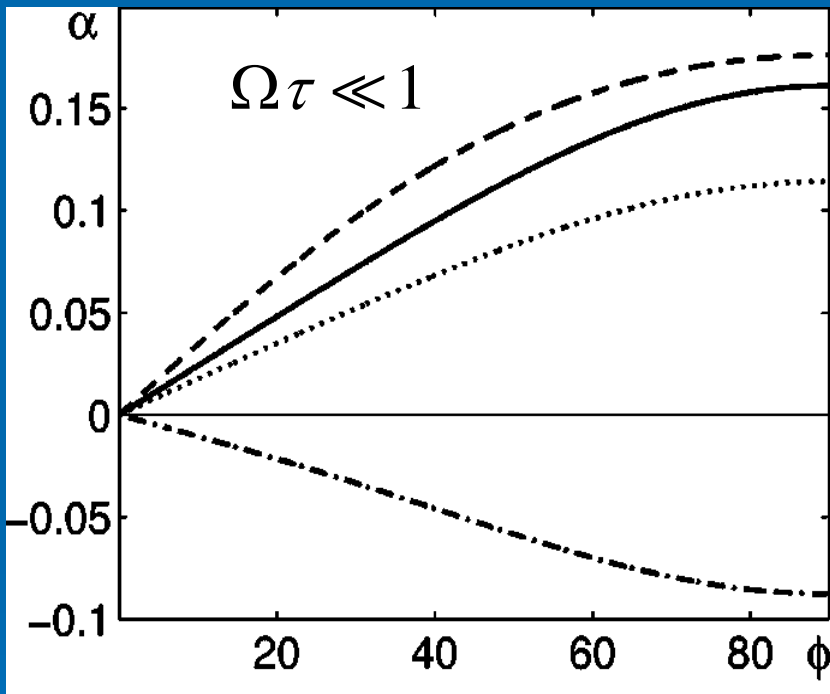
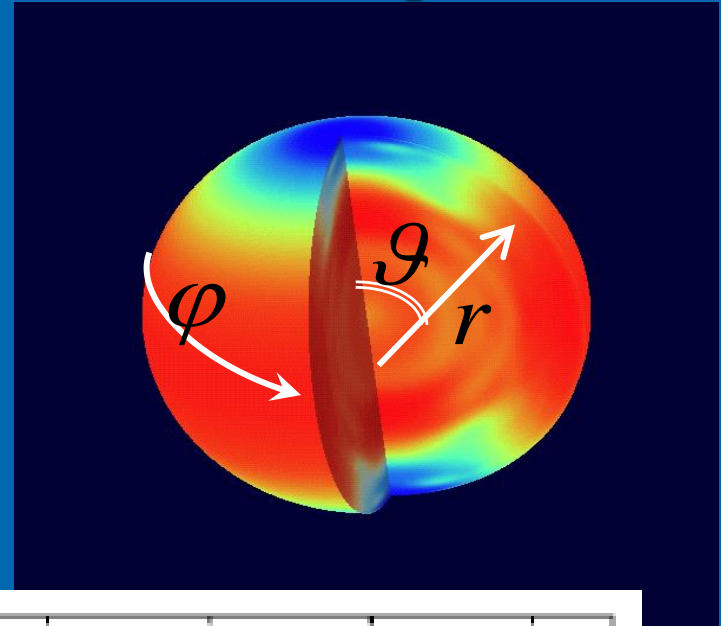
Kichatinov, L. L. & Rudiger, G. (1993)

Kleorin & Rogachevskii (2003)

Latitudinal profiles of the alpha-effect

$$\phi = \frac{\pi}{2} - \vartheta$$

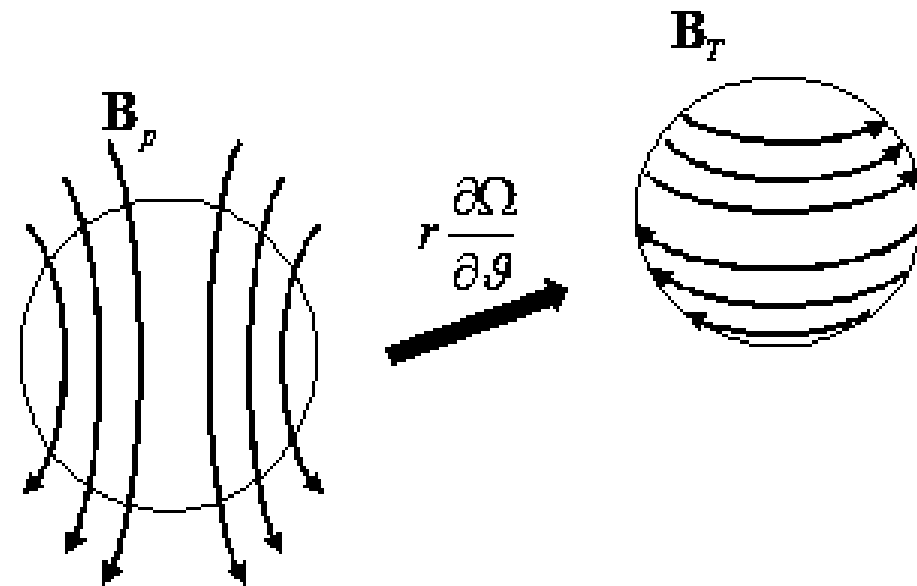
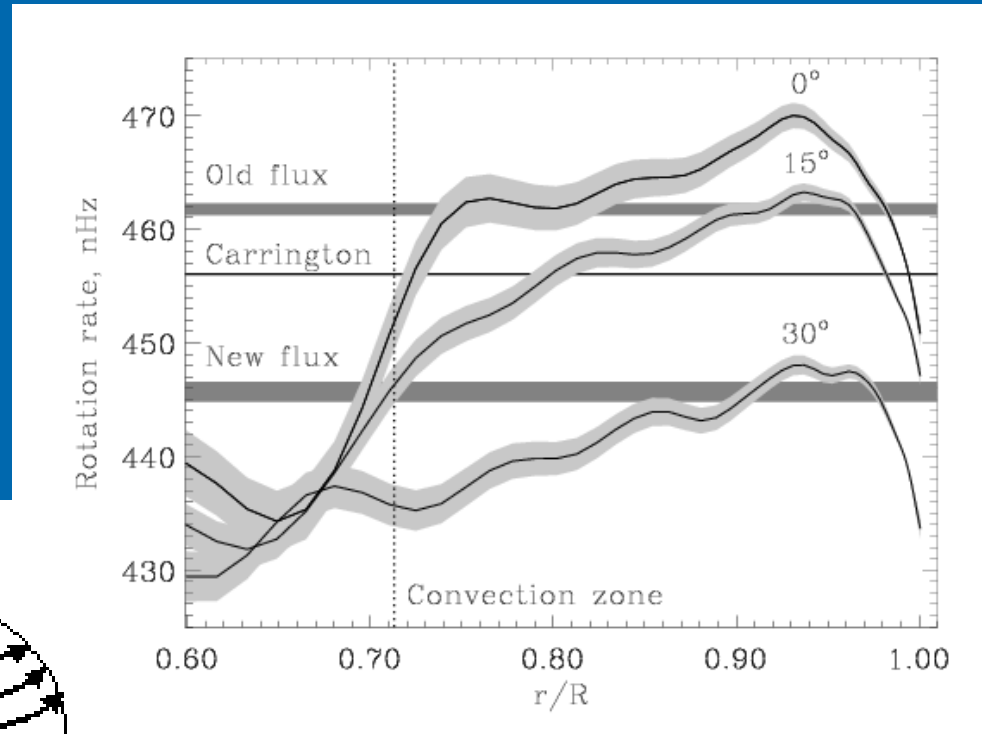
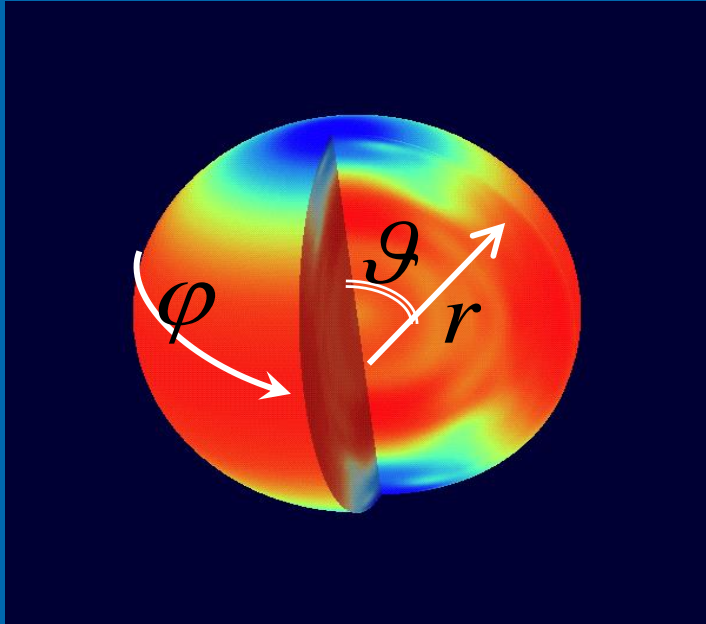
effect



Kleorin & Rogachevskii (2003)

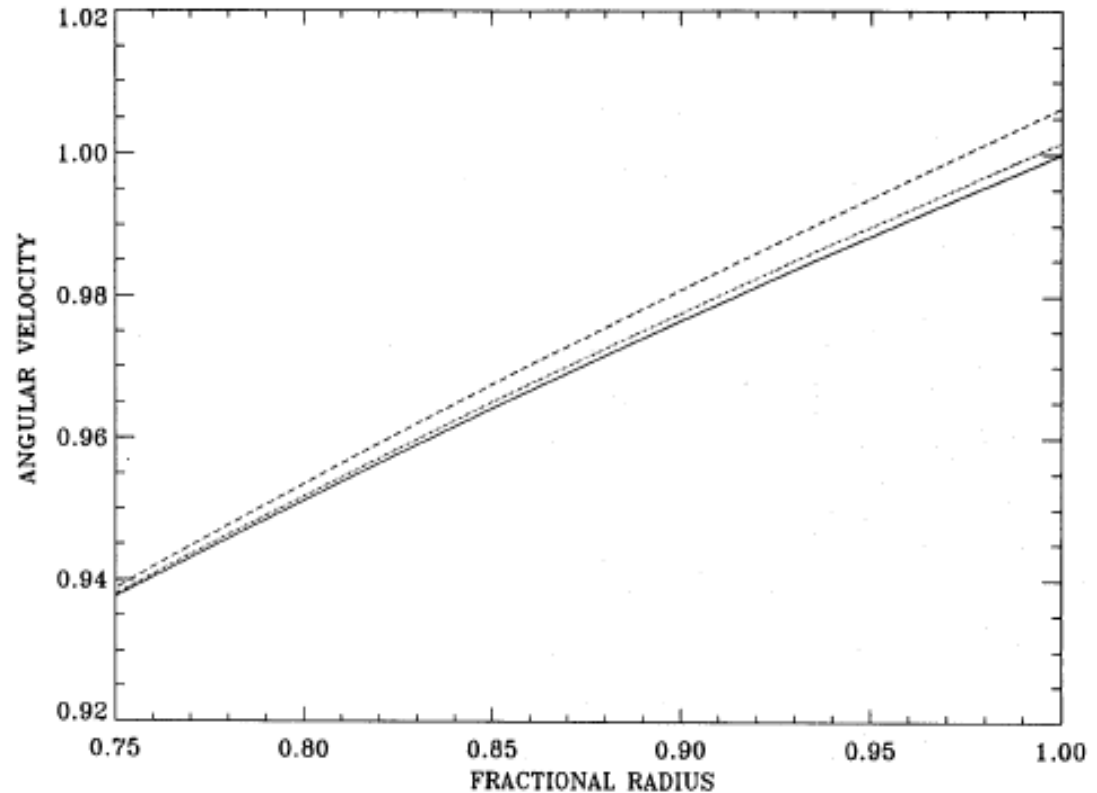
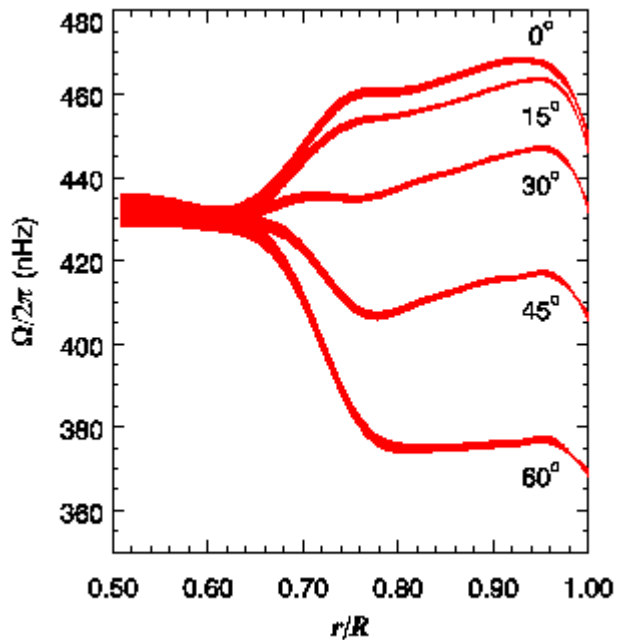
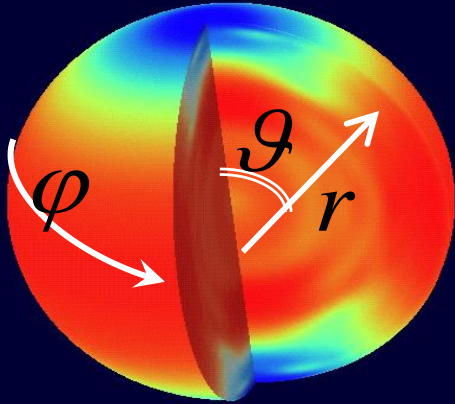
Differential Rotation

$$\mathbf{U} = \mathbf{e}_\varphi \Omega(r, \vartheta) r \sin(\vartheta)$$



Differential Rotation

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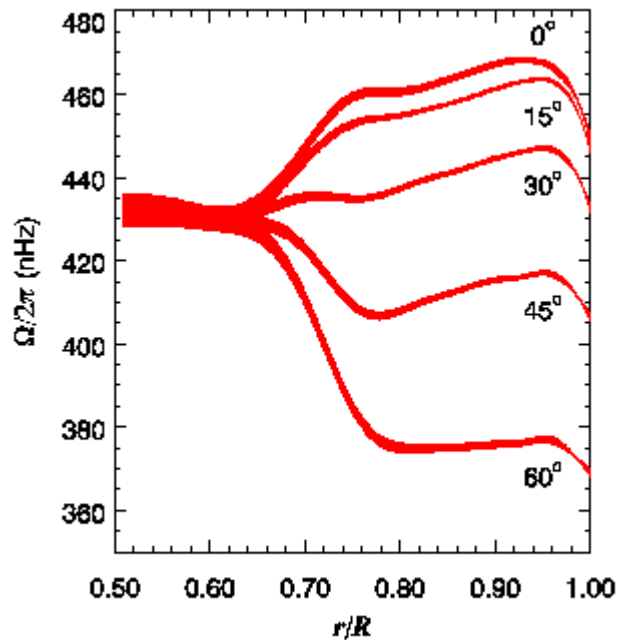
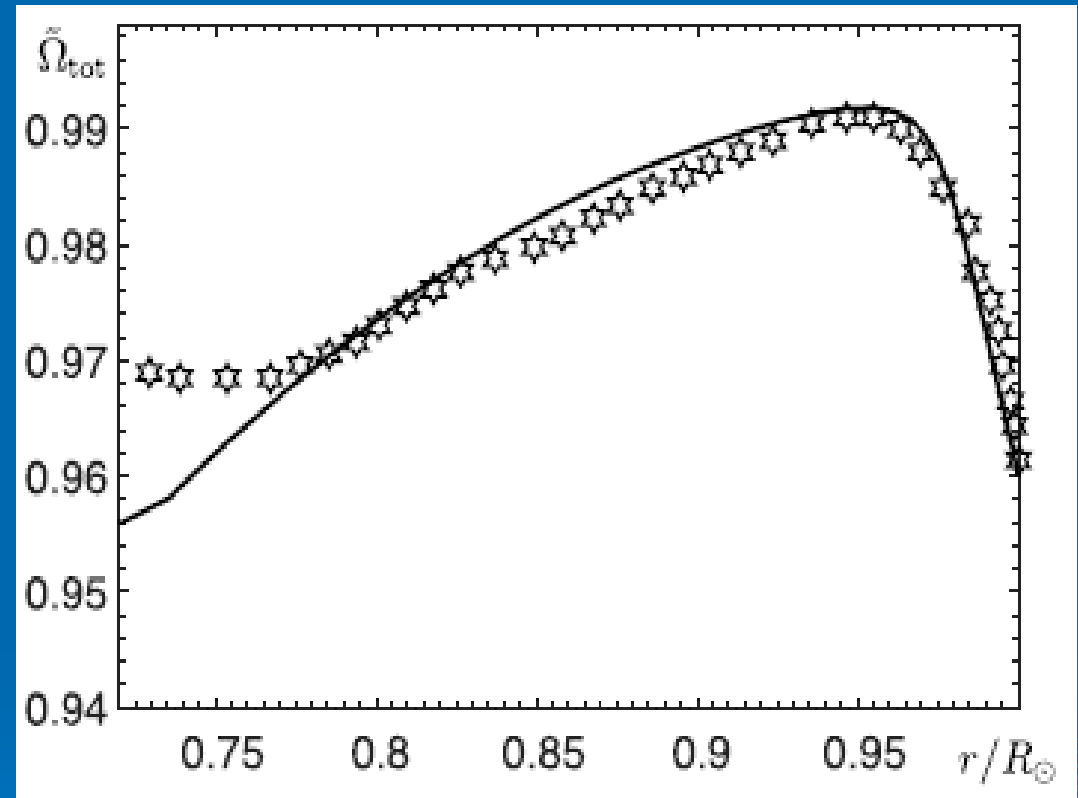
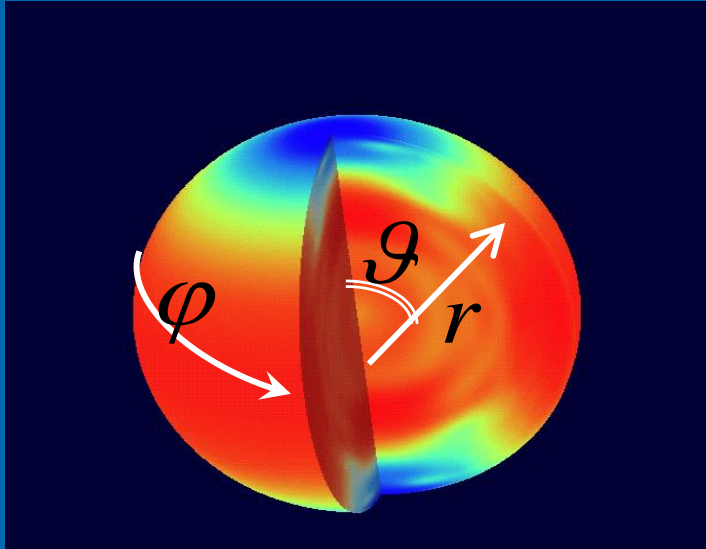


Kichatinov, L. L. & Rudiger, G. (1993)

Differential Rotation

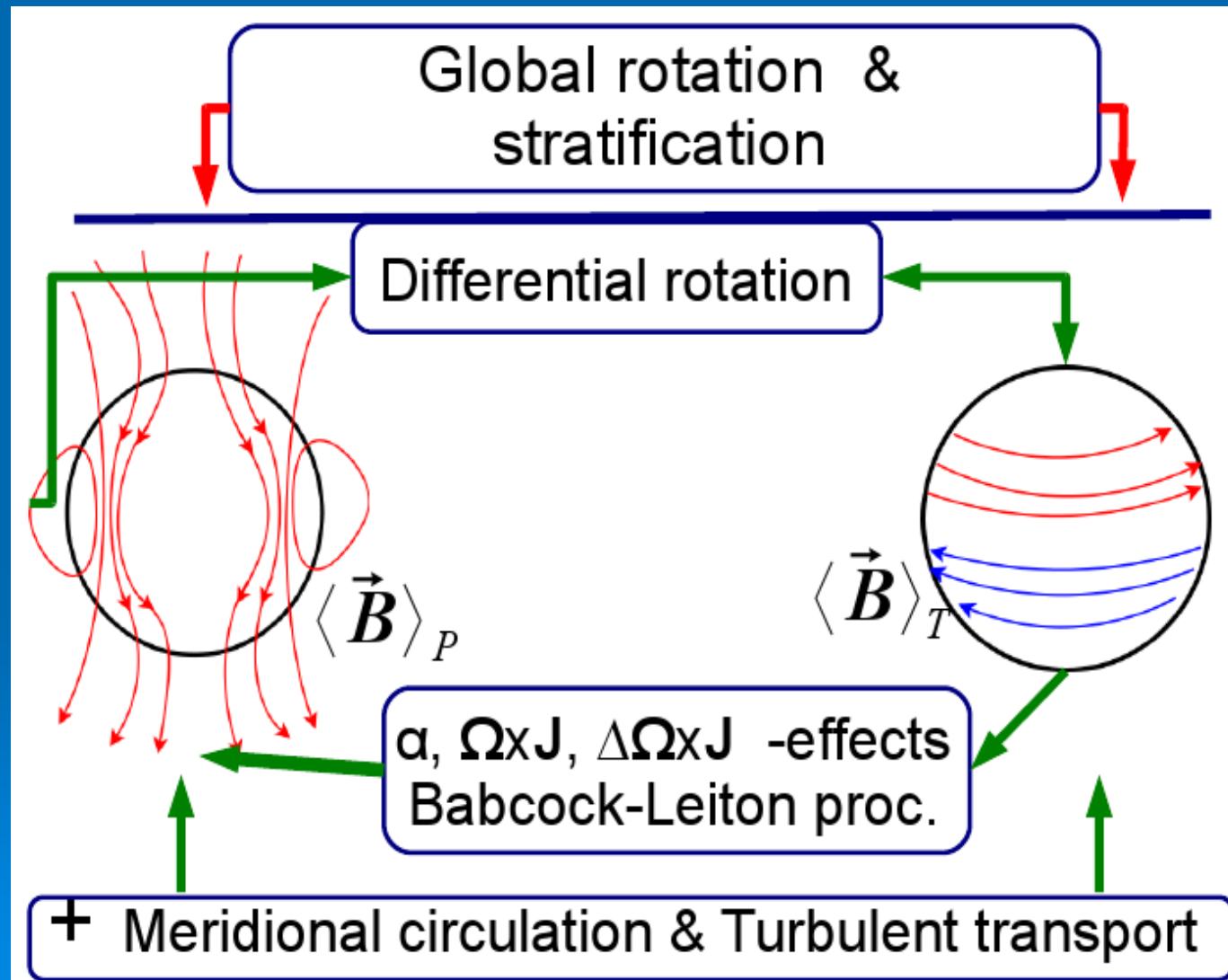
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$$\mathbf{U} = \mathbf{e}_\varphi \Omega(r, \vartheta) r \sin(\vartheta)$$



Rogachevskii & Kleeorin (2018)

Basic Mechanisms

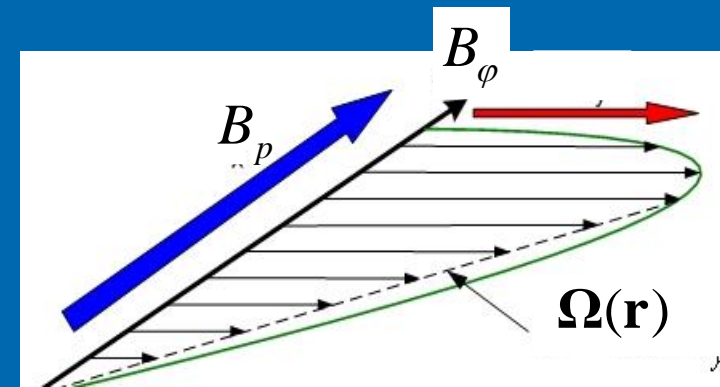


Generation of the mean magnetic field due to the $\alpha\Omega$ dynamo

Mean magnetic field:

$$\mathbf{B} = \mathbf{e}_\varphi B(r, \vartheta) + \text{rot} \left[\mathbf{e}_\varphi A(r, \vartheta) \right]$$

$$r \sin(\vartheta) [\nabla \Omega \times \nabla A]_\varphi = \frac{1}{r} \left(\frac{\partial \Omega}{\partial r} \frac{\partial}{\partial \vartheta} - \frac{\partial \Omega}{\partial \vartheta} \frac{\partial}{\partial r} \right) (rA \sin(\vartheta))$$

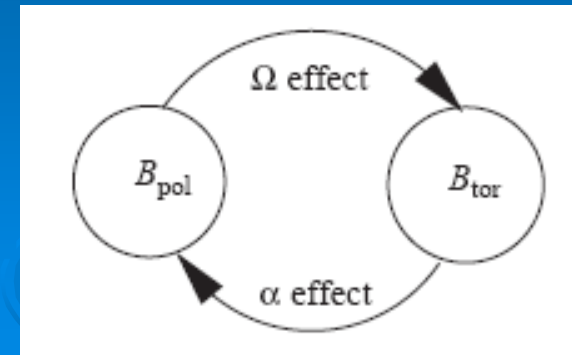


$$\frac{\partial B}{\partial t} = rD \sin(\vartheta) [\nabla \Omega \times \nabla A]_\varphi + \eta_T \Delta_s B - R_\alpha^2 \hat{\Delta}(\alpha A)$$

$$\frac{\partial A}{\partial t} = \alpha B + \Delta_s A; \quad \Delta_s = \Delta - (r \sin \vartheta)^{-2}$$

Dynamo number:

$$D = R_\alpha R_\Omega; \quad R_\alpha = \frac{\alpha R_\odot}{\eta_T}; \quad R_\Omega = \frac{\delta \Omega R_\odot^2}{\eta_T};$$



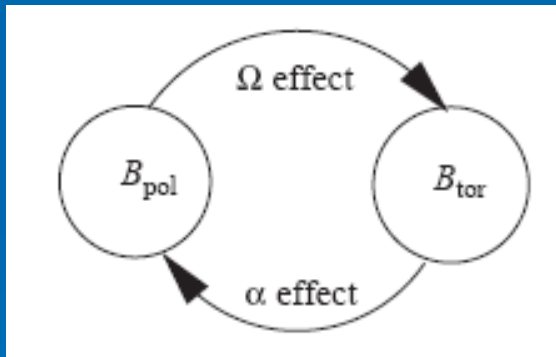
Generation of the mean magnetic field due to the $\alpha\Omega$ dynamo

Mean magnetic field:

$$\mathbf{B} = B_\varphi \mathbf{e}_\varphi + \nabla \times [A \mathbf{e}_\varphi]$$

$$\left(\frac{\partial}{\partial t} - \Delta \right) B = D \sin(\vartheta_0) \frac{\partial \Omega}{\partial r} \frac{\partial A}{\partial \vartheta} \left| \left(\frac{\partial}{\partial t} - \eta_T \Delta \right) \right.$$

$$\left. \left(\frac{\partial}{\partial t} - \Delta \right) A = \cos(\vartheta_0) B \right.$$



Dynamo number:

$$D = R_\alpha R_\Omega; \quad R_\alpha = \frac{\alpha_0 R_\odot}{\eta_T}; \quad R_\Omega = \frac{\delta \Omega R_\odot^2}{\eta_T};$$

$$\Delta_s = \Delta - \frac{1}{r^2 \sin^2 \theta}$$

Generation of the mean magnetic field due to the $\alpha\Omega$ dynamo

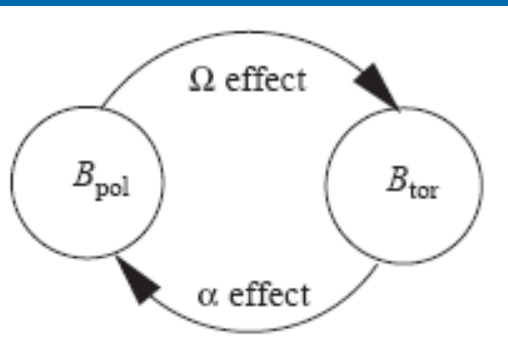
Mean magnetic field:

$$\mathbf{B} = B_\varphi \mathbf{e}_\varphi + \nabla \times [A \mathbf{e}_\varphi]$$

$$B = B_0 \exp\left(\lambda t - i(K_r r + K_g \vartheta)\right)$$

$$\left(\frac{\partial}{\partial t} - \Delta\right)^2 B = D \cos(\vartheta_0) \sin(\vartheta_0) \frac{\partial \Omega}{\partial r} \frac{\partial B}{\partial \vartheta}$$

$$\left(\frac{\partial}{\partial t} - \Delta\right) A = \cos(\vartheta_0) B$$



Dynamo number:

$$\Delta_s = \Delta - \frac{1}{r^2 \sin^2 \theta}$$

$$D = R_\alpha R_\Omega; \quad R_\alpha = \frac{\alpha_0 R_\odot}{\eta_T}; \quad R_\Omega = \frac{\delta \Omega R_\odot^2}{\eta_T};$$

Generation of the mean magnetic field due to the $\alpha\Omega$ dynamo

Mean magnetic field: $B = B_0 \exp\left(\lambda t - i(K_r r + K_g \mathcal{G})\right) \leftarrow$ E. Parker (1955)

$$\mathbf{B} = B_\varphi \mathbf{e}_\varphi + \nabla \times [A \mathbf{e}_\varphi]$$

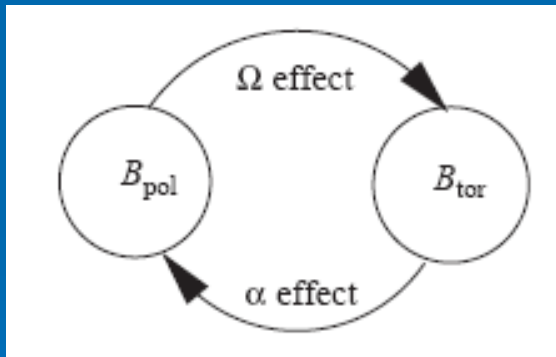
$$\left(\lambda + (K_r^2 + K_g^2)\right)^2 B_0 = -B_0 \frac{iK_g}{2} D \sin(2\mathcal{G}_0)$$

$$\left(\frac{\partial}{\partial t} - \Delta\right) A = \cos(\mathcal{G}_0) B$$

Dynamo number:

$$D = R_\alpha R_\Omega; \quad R_\alpha = \frac{\alpha_0 R_\odot}{\eta_T}; \quad R_\Omega = \frac{\delta \Omega R_\odot^2}{\eta_T};$$

$$\Delta_s = \Delta - \frac{1}{r^2 \sin^2 \theta}$$



Generation of the mean magnetic field due to the $\alpha\Omega$ dynamo

Mean magnetic field:

$$\mathbf{B} = B_\varphi \mathbf{e}_\varphi + \nabla \times [A \mathbf{e}_\varphi]$$

$$B = B_0 \exp(\gamma t) \cos(\omega t - i(K_r r + K_g \vartheta))$$

$$\gamma = \frac{1}{2} \sqrt{|K_g \sin(2\vartheta_0) D|} - (K_r^2 + K_g^2); \quad \frac{d\gamma}{dK_g} = 0$$

Dynamo number:

$$K_0 = \frac{1}{2} \sqrt[3]{\frac{|\sin(2\vartheta_0) D|}{2}};$$

$$D = R_\alpha R_\Omega; \quad R_\alpha = \frac{\alpha_0 R_\odot}{\eta_T}; \quad R_\Omega = \frac{\delta \Omega R_\odot^2}{\eta_T};$$

$$\gamma(K_0) = \left[\frac{1}{2} \left(\frac{|\sin(2\vartheta_0) D|}{2} \right)^{2/3} - \frac{1}{4} \left(\frac{|\sin(2\vartheta_0) D|}{2} \right)^{2/3} - K_r^2 \right]; \quad D = \alpha_0 \frac{\partial \Omega}{\partial r} \frac{R_\odot^4}{\eta_T^2} < 0$$

Generation of the mean magnetic field due to the $\alpha\Omega$ dynamo

$$B = B_0 \exp(\gamma t) \cos(\gamma t - K_0 \vartheta + \delta) \cos(K_r (1 - r))$$

Mean magnetic field:

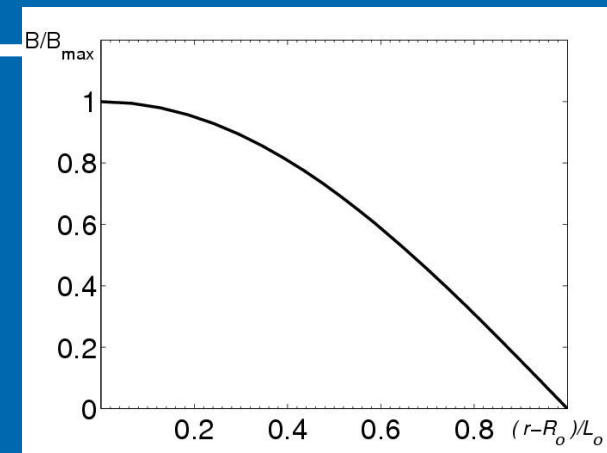
$$\mathbf{B} = B_\varphi \mathbf{e}_\varphi + \nabla \times [A \mathbf{e}_\varphi]$$

$$K_0 = \frac{1}{2} \sqrt[3]{\frac{|\sin(2\vartheta_0) D|}{2}};$$

$$K_r^{(\min)} = R_\odot \frac{\pi}{2L_\odot}$$

$$\gamma(K_0) = \left[\frac{1}{4} \left(\frac{|\sin(2\vartheta_0) D|}{2} \right)^{2/3} - R_\odot^2 K_r^2 \right] < \left[\frac{1}{4} \left(\frac{|D|}{2} \right)^{2/3} - K_r^2 \right]$$

$$D = \alpha_0 \frac{\partial \Omega}{\partial r} \frac{R_\odot^4}{\eta_T^2} < 0; |D| > D_{cr} = 2(2K_r)^3 \approx 2(\pi R_\odot / L_\odot)^3 \approx \underline{1,680}$$

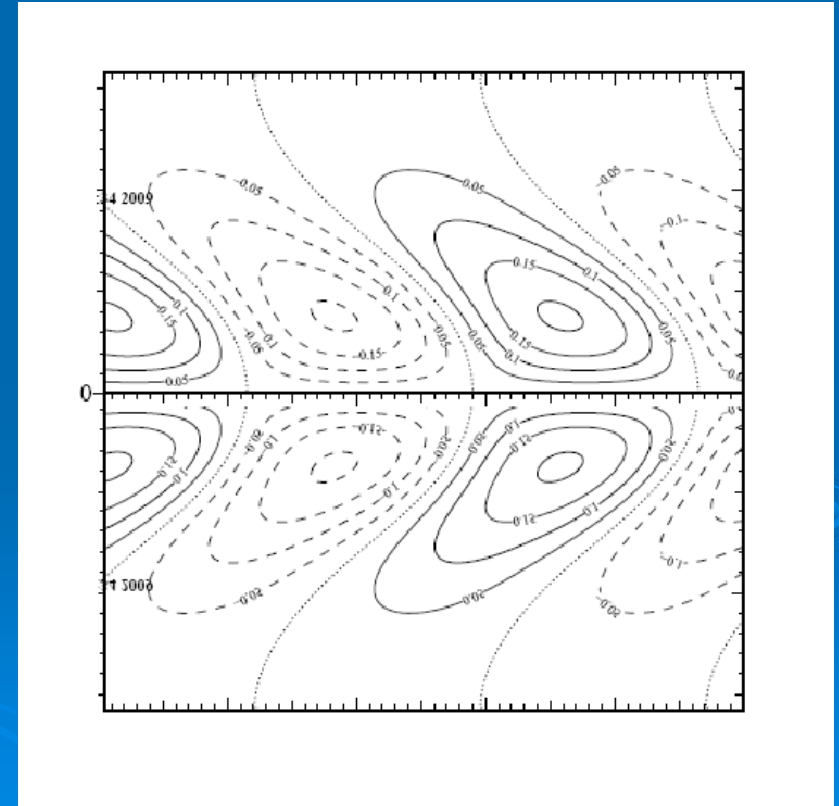
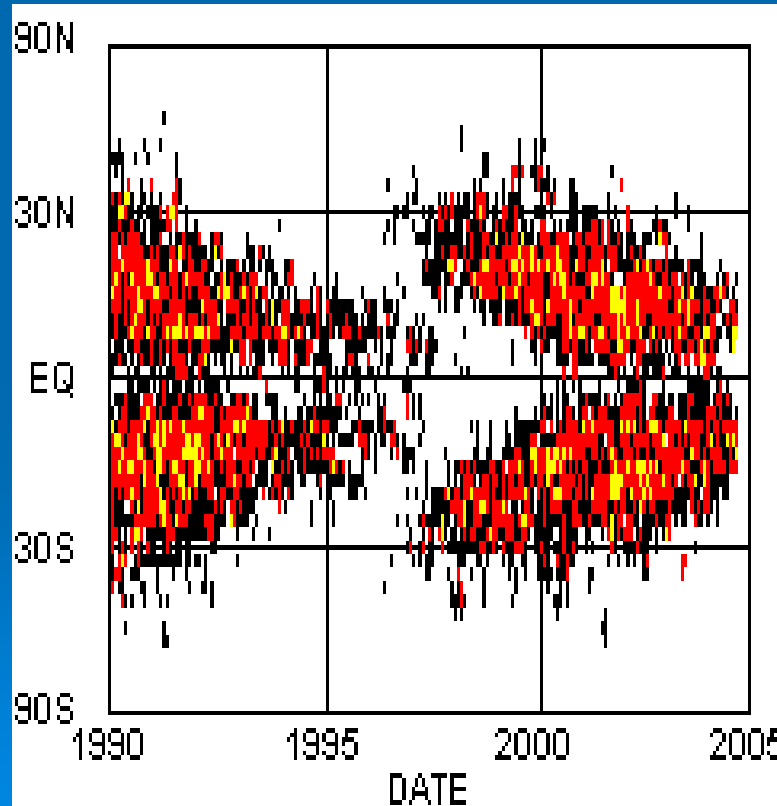


Comparing Observations and Theory

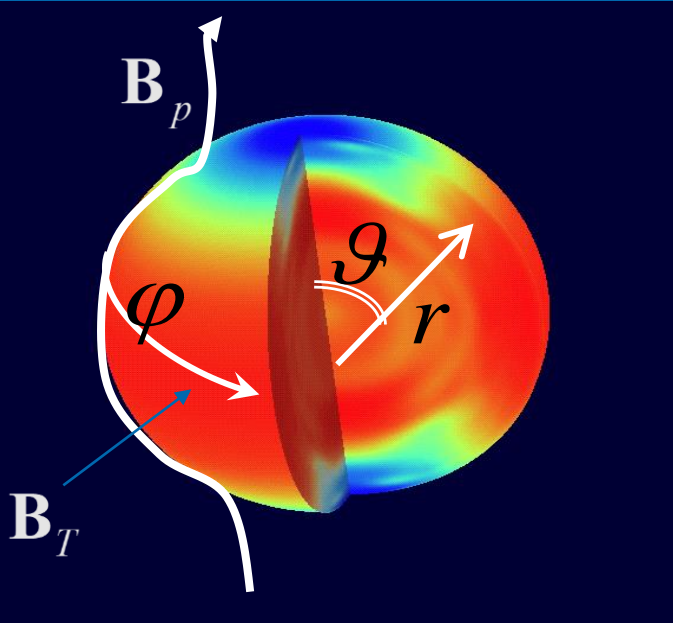
$B_{00} - ?$

$$B = B_{00} \exp(\gamma t) \cos(\gamma t - K_0 \mathcal{G} + \delta) \cos(K_r (1 - r))$$

Steenbeck, Krause and Raedler (1966)



Nonlinear Alpha-Omega Dynamo (Mean-Field Approach)



$$\frac{\partial}{\partial \varphi} = 0 \quad \mathbf{U} = \mathbf{e}_\varphi \Omega(r, \vartheta) r \sin(\vartheta);$$

Iroshnikov (1970)

$$\alpha(r, \vartheta, B) = \frac{\alpha(r, \vartheta)}{1 + (B/B_0)^2}$$

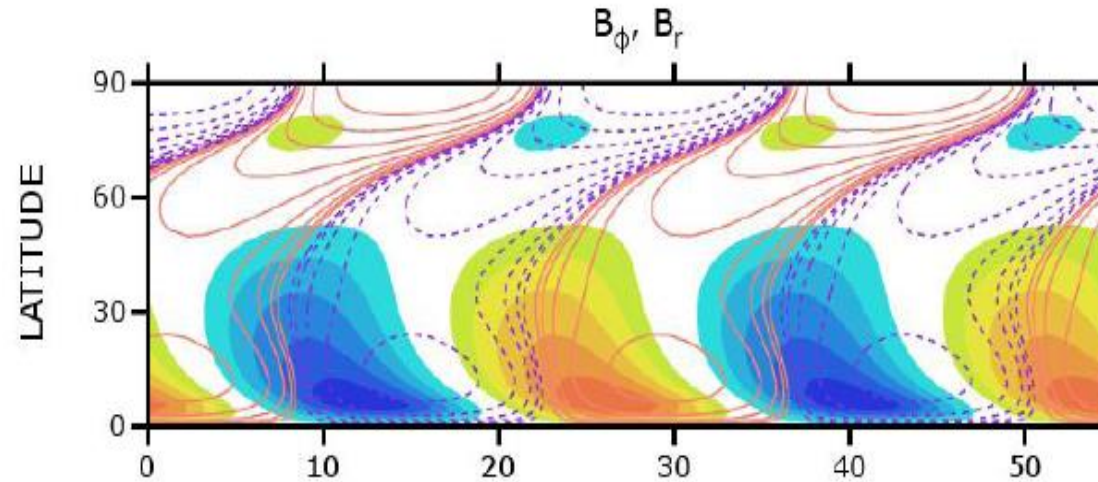
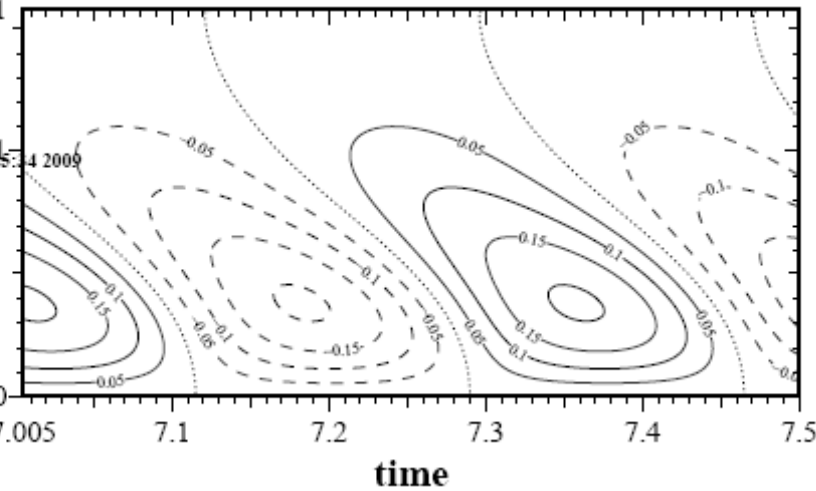
$$\frac{\partial B}{\partial t} = (\mathbf{B}_p \cdot \nabla) \mathbf{U}_T + \eta_T \Delta_s B$$

$$\frac{\partial A}{\partial t} = \alpha(r, \vartheta, B) B + \eta_T \Delta_s A; \quad \Delta_s = \Delta - (r \sin \vartheta)^{-2}$$

$B_0 - ?$ Nonlinear Alpha-Omega Dynamo (Mean-Field Approach)

Steenbeck, Krause and Raedler (1966)

~ e.g., Pipin (2000)



$$\frac{\partial B}{\partial t} = (\mathbf{B}_p \cdot \nabla) \mathbf{U}_T + \eta_T \Delta_s B$$

$$\alpha(r, \vartheta, B) = \frac{\alpha(r, \vartheta)}{1 + (B/B_0)^2}$$

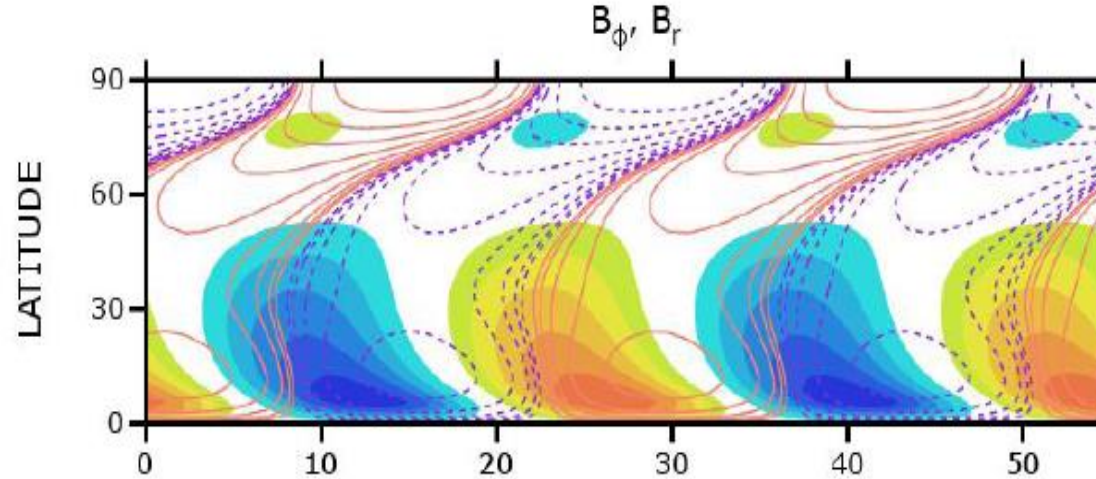
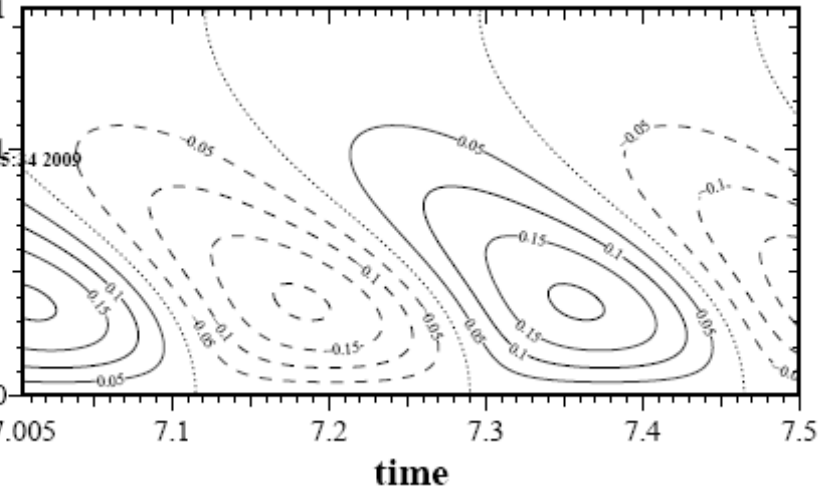
Iroshnikov (1970)

$$\frac{\partial A}{\partial t} = \alpha(r, \vartheta, B) B + \eta_T \Delta_s A; \quad \Delta_s = \Delta - (r \sin \vartheta)^{-2}$$

B_{00} – ? Nonlinear Alpha-Omega Dynamo (Mean-Field Approach)

Steenbeck, Krause and Raedler (1966)

Pipin (2000)



$$\frac{\partial B}{\partial t} = (\mathbf{B}_p \cdot \nabla) \mathbf{U}_T + \eta_T \Delta_s B \rightarrow$$

$$\rightarrow D(B) = \frac{D}{1 + (B/B_0)^2} \rightarrow D_{cr}$$

$$\frac{\partial A}{\partial t} = \alpha(r, \vartheta, B) B + \eta_T \Delta_s A;$$

$$\Delta_s = \Delta - (r \sin \vartheta)^{-2}$$

Iroshnikov (1970)

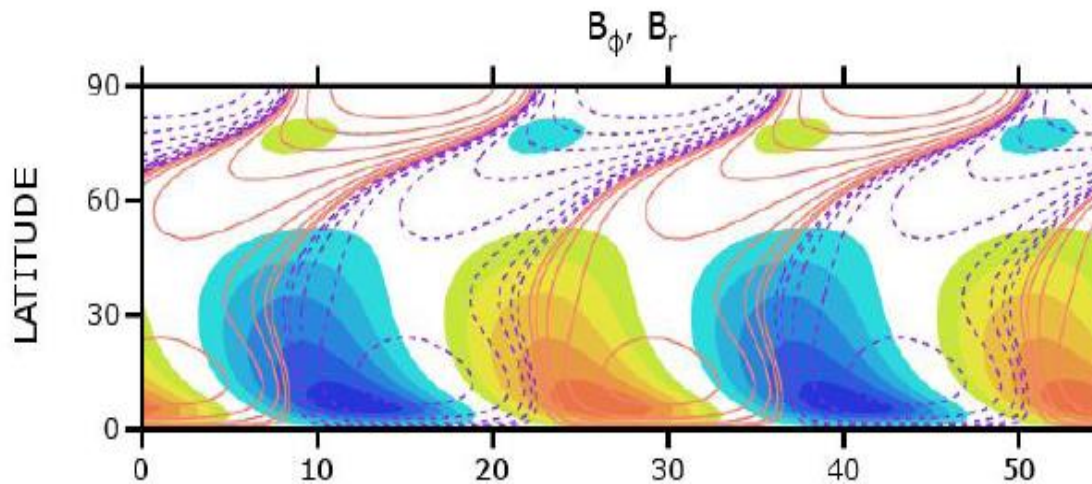
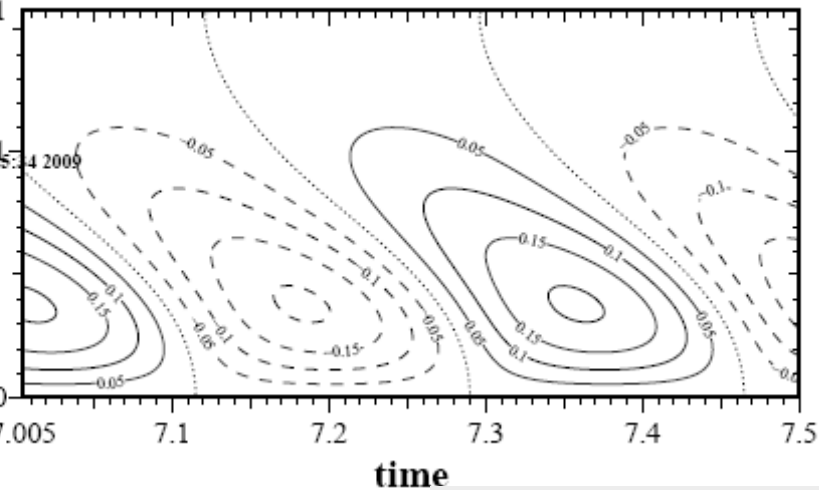
$$\alpha(r, \vartheta, B) = \frac{\alpha(r, \vartheta)}{1 + (B/B_0)^2} \rightarrow$$

$$\rightarrow B \approx B_0 \sqrt{\frac{D}{D_{cr}} - 1}$$

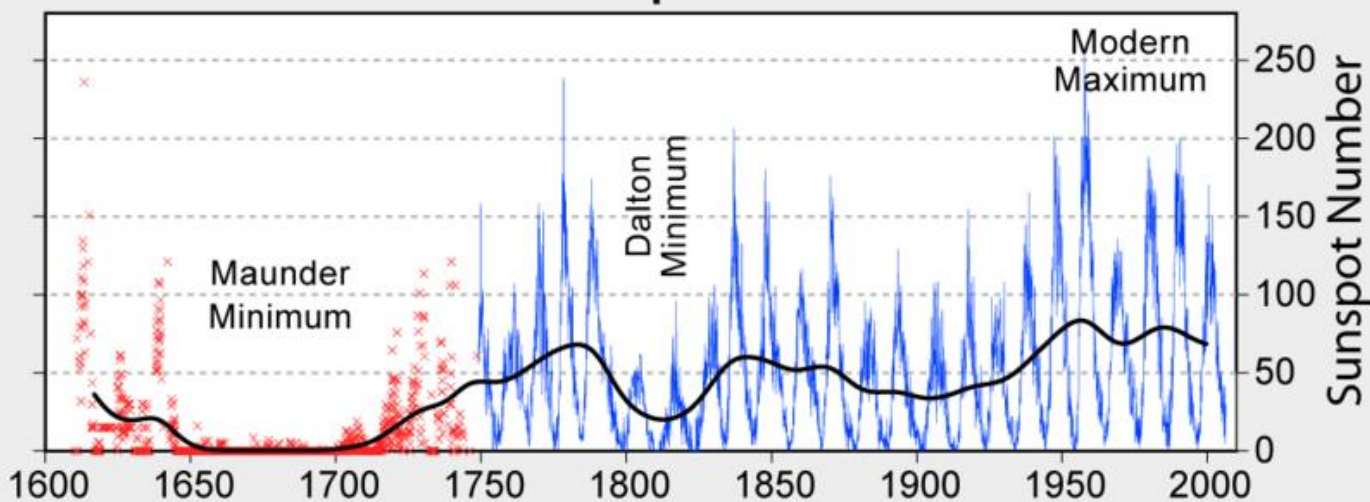
B_{00} – ? Nonlinear Alpha-Omega Dynamo (Mean-Field Approach)

Steenbeck, Krause and Raedler (1966)

Pipin (2000)



400 Years of Sunspot Observations

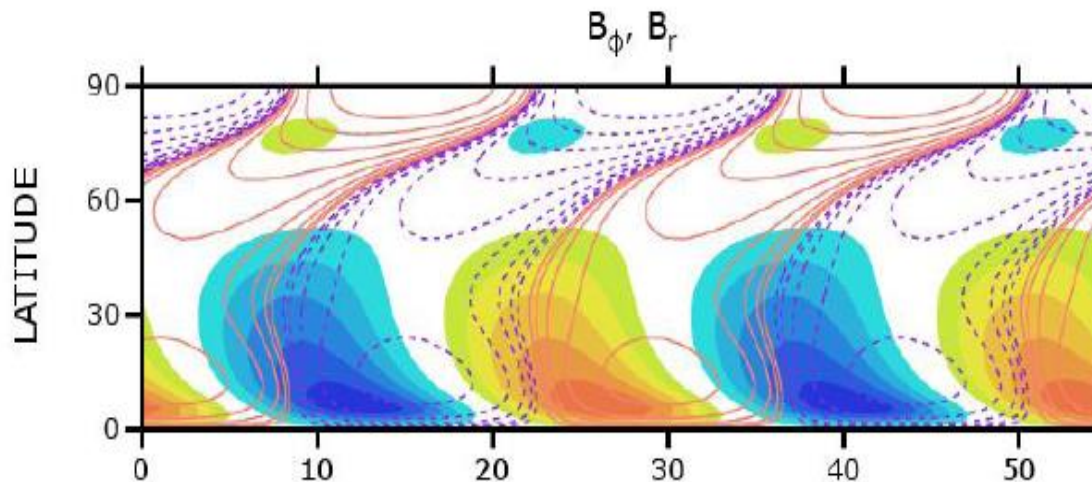
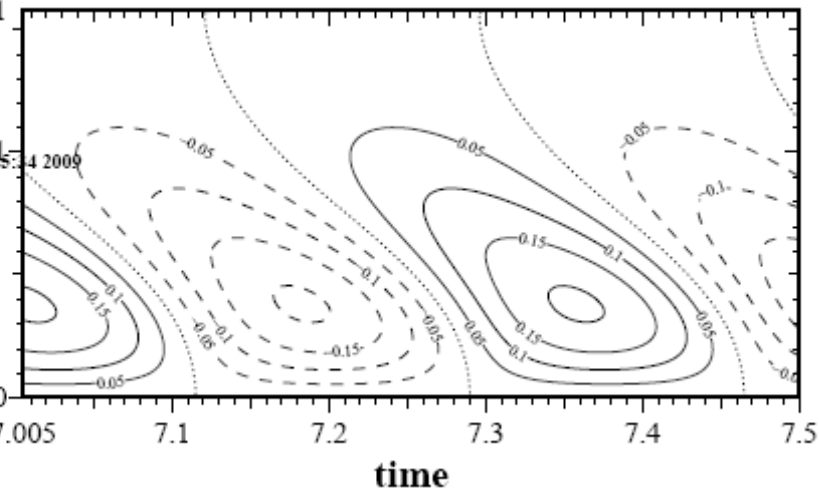


$$W \approx 11N_{Sp}$$

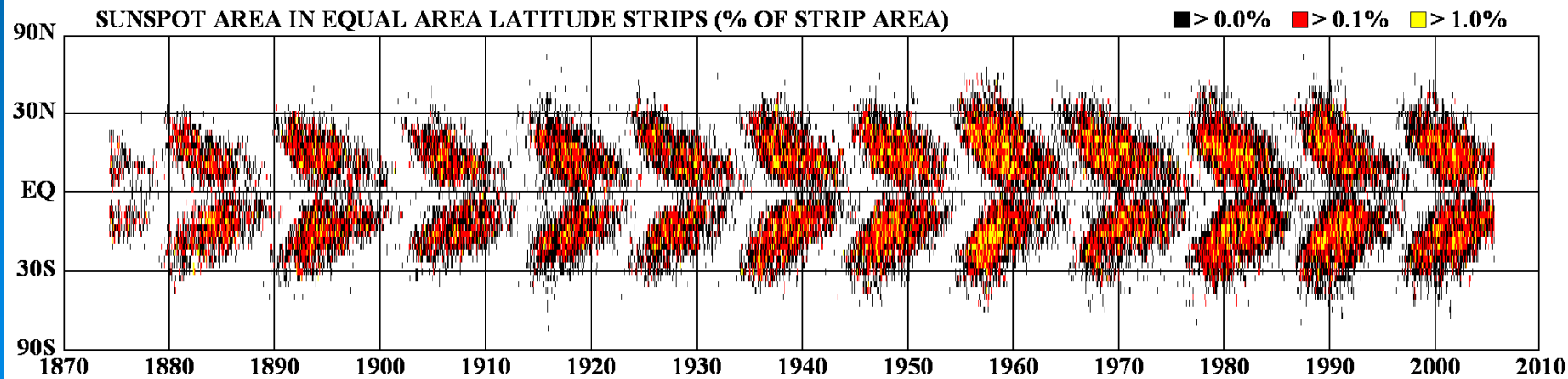
B_{00} – ? Nonlinear Alpha-Omega Dynamo (Mean-Field Approach)

Steenbek, Krause and Redler (1966)

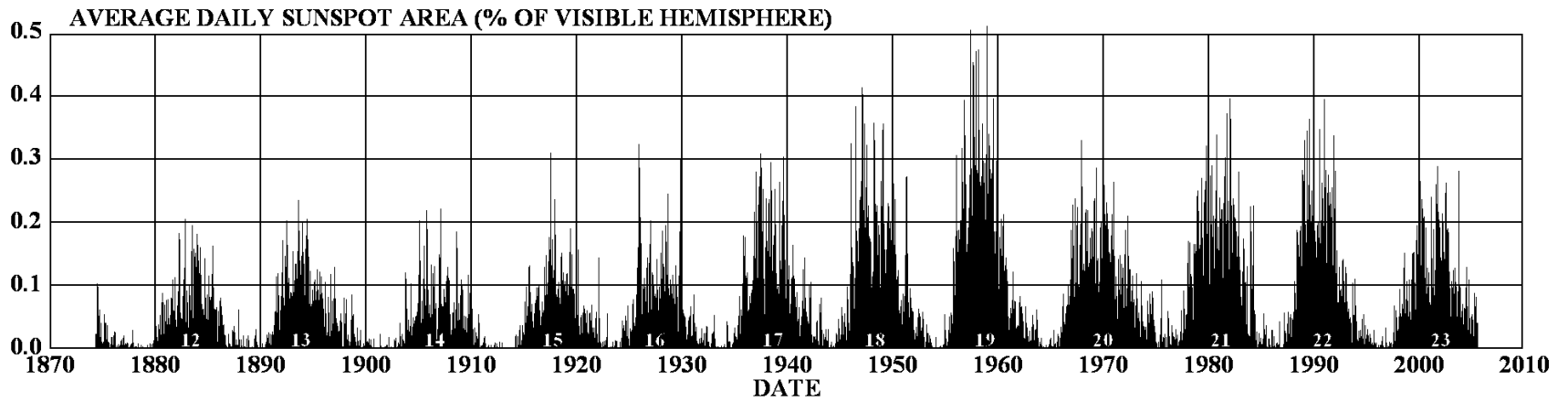
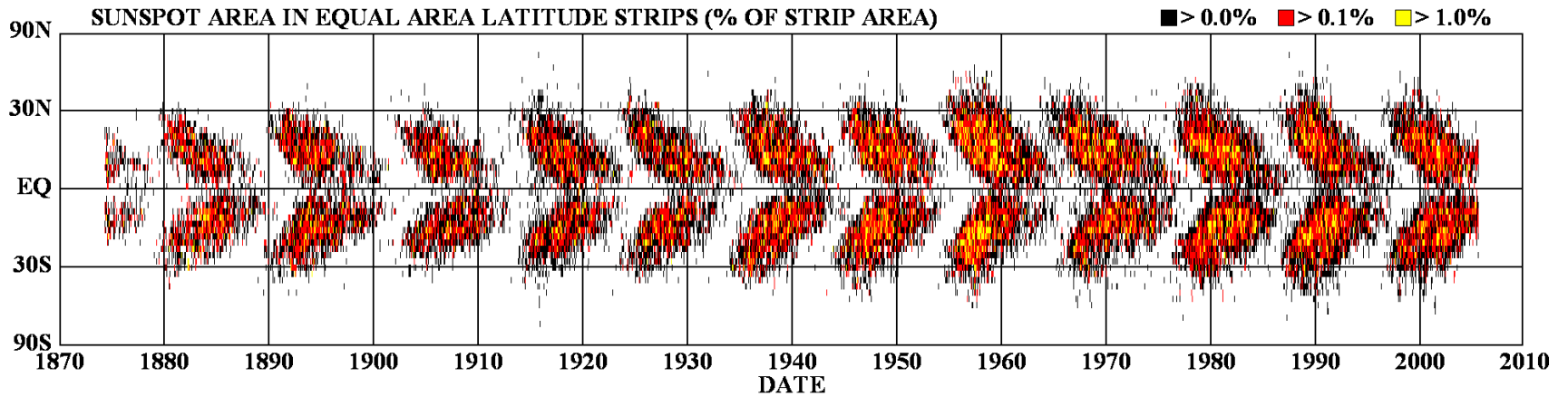
Pipin (2000)



DAILY SUNSPOT AREA AVERAGED OVER INDIVIDUAL SOLAR ROTATIONS



DAILY SUNSPOT AREA AVERAGED OVER INDIVIDUAL SOLAR ROTATIONS



Nonlinear Dynamo



Magnetic part of Alpha -Effect (Mean-Field Approach)

- Induction equation for **mean magnetic field**:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{U} \times \mathbf{B} + \boldsymbol{\varepsilon} - \eta \nabla \times \mathbf{B})$$

- **Electromotive force**:

$$\boldsymbol{\varepsilon} \equiv \langle \mathbf{u} \times \mathbf{b} \rangle = \boldsymbol{\alpha} \mathbf{B} - \eta_T \nabla \times \mathbf{B} + \dots$$

$$\boldsymbol{\alpha} = -\frac{\tau}{3} \langle \mathbf{u} \cdot \text{rot } \mathbf{u} \rangle + \frac{\tau}{12\pi\rho} \underbrace{\langle \mathbf{b} \cdot \text{rot } \mathbf{b} \rangle}_{\sim \mathbf{a} \cdot \mathbf{b}}$$

Evolution equation magnetic for part of Alpha -Effect

$$\frac{\partial \alpha_m}{\partial t} + \text{div } \mathbf{F} = \frac{Q}{4\pi\rho\eta_T} \left[-(\alpha_h + \alpha_m) \mathbf{B}^2 + \eta_T \mathbf{B} \cdot \text{rot} \mathbf{B} \right] - \frac{\alpha_m}{\tau_\chi}$$

$$\tau_\chi \approx \tau_0 \text{Rm} \quad \text{Kleeorin \& Ruzmaikin (1982)}$$

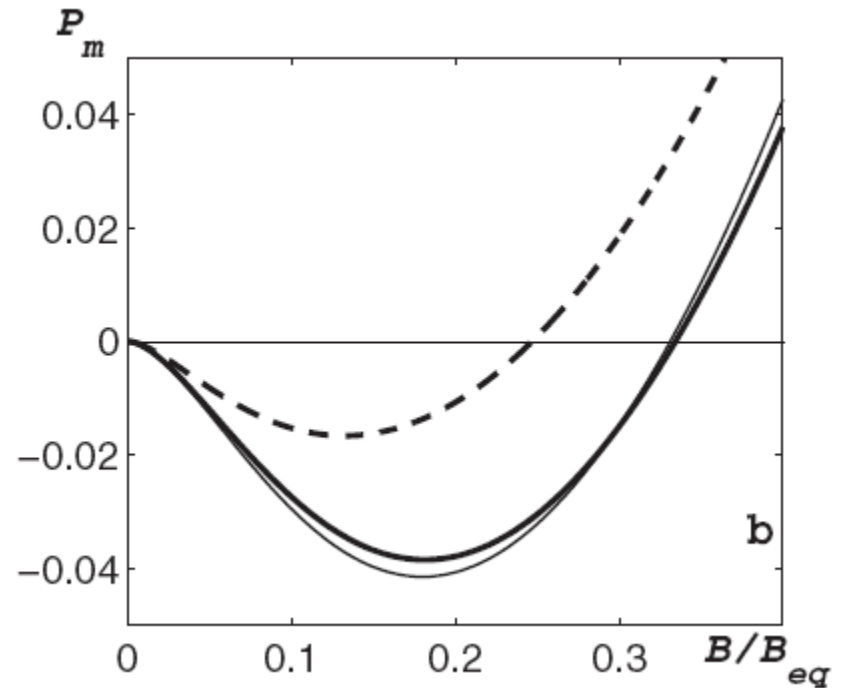
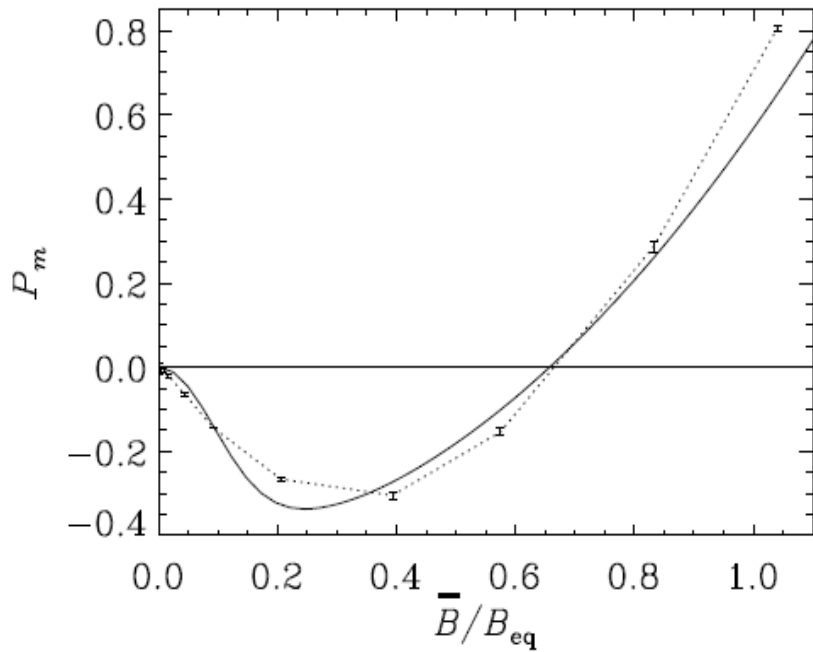
➤ Electromotive force:

$$\boldsymbol{\varepsilon} \equiv \langle \mathbf{u} \times \mathbf{b} \rangle = \boldsymbol{\alpha} \mathbf{B} - \boldsymbol{\eta}_T \nabla \times \mathbf{B} + \dots$$

$$\alpha = \underbrace{-\frac{\tau}{3} \langle \mathbf{u} \cdot \text{rot } \mathbf{u} \rangle}_{\alpha_h} \varphi_k(B) + \underbrace{\frac{\tau}{12\pi\rho} \langle \mathbf{b} \cdot \text{rot } \mathbf{b} \rangle}_{\alpha_m, \sim \mathbf{a} \cdot \mathbf{b}} \varphi_m(B)$$

Field & Blackman (2000)

Negative Magnetic Pressure



(Brandenburg, Kleeorin, Rogachevskii 2010)

$$P_m = (1 - q_p) \frac{\overline{B}^2}{B_{eq}^2}$$

(Kleeorin, Rogachvskii, Ruzmaikin
1989,1990; Rogachvskii, Kleeorin
2007)

Negative Magnetic Pressure

МОСКОВСКИЙ ОРДЕНА ЛЕНИНА И ОРДЕНА ТРУДОВОГО КРАСНОГО ЗНАМЕНИ
ГОСУДАРСТВЕННЫЙ ПЕДАГОГИЧЕСКИЙ ИНСТИТУТ имени В.И.ЛЕНИНА

На правах рукописи

КЛИОРИН НАТАН ИОСИФОВИЧ

УДК 524.3-337+523.9-337

НЕЛИНЕЙНОЕ ДИНАМО КРУПНОМАСШТАБНЫХ МАГНИТНЫХ
ПОЛЕЙ ЗВЕЗД

(Специальность - 01.03.02 - астрофизика)

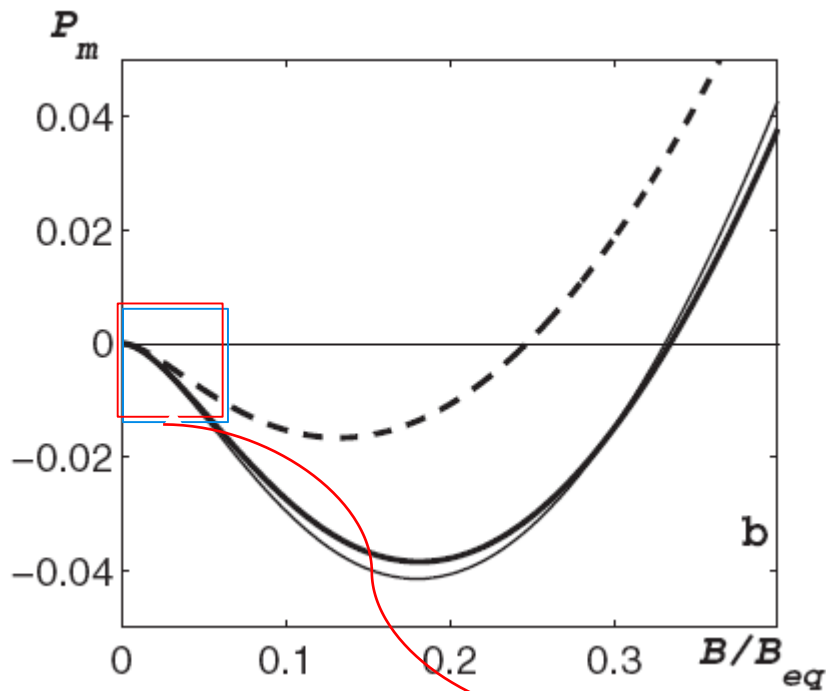
Диссертация на соискание
ученой степени кандидата
физико-математических наук



Научный руководитель:
доктор физико-математических
наук, старший научный сотрудник
А.А.РУЗМАЙКИН

Москва - 1984

Negative Magnetic Pressure



где $K_1 \approx (1 - 2/15 \ln R_m)$,

$K = (1 - 2/45 \ln R_m)$.

$$\sigma_{ij}^M = \frac{1}{4\pi} \left[-\frac{K_1}{2} B^2 \delta_{ij} + K B_i B_j + \frac{\langle h_0^2 \rangle}{6} \delta_{ij} \right],$$

(I.3.II)

How to Understand Negative Magnetic Pressure?

$$P_B = \frac{\bar{B}^2}{8\pi}; \quad p_T = \frac{1}{3}E_M + \frac{2}{3}E_K$$

$$p_T = -\frac{1}{24\pi}\langle b^2 \rangle + \frac{2}{3}E_T^{(tot)}; \quad E_T^{(tot)} = \frac{\rho_0 \langle u^2 \rangle}{2} + \frac{\langle b^2 \rangle}{8\pi}$$

$$P_{MHD} = P^{(tot)} - P_g = \frac{1}{8\pi} \left(\bar{B}^2 - \frac{\langle b^2 \rangle}{3} \right) + \frac{2}{3}E_T^{(tot)};$$

$$\langle b^2 \rangle = a\bar{B}^2; \quad E_T^{(tot)} = E_T^{(tot)}(0) - \frac{a_1}{8\pi}\bar{B}^2;$$

$$P_{MHD} = \frac{\bar{B}^2}{8\pi} \left(1 - \frac{a + 2a_1}{3} \right) + \frac{2}{3}E_T^{(tot)}(0)$$

How to Understand Negative Magnetic Pressure?

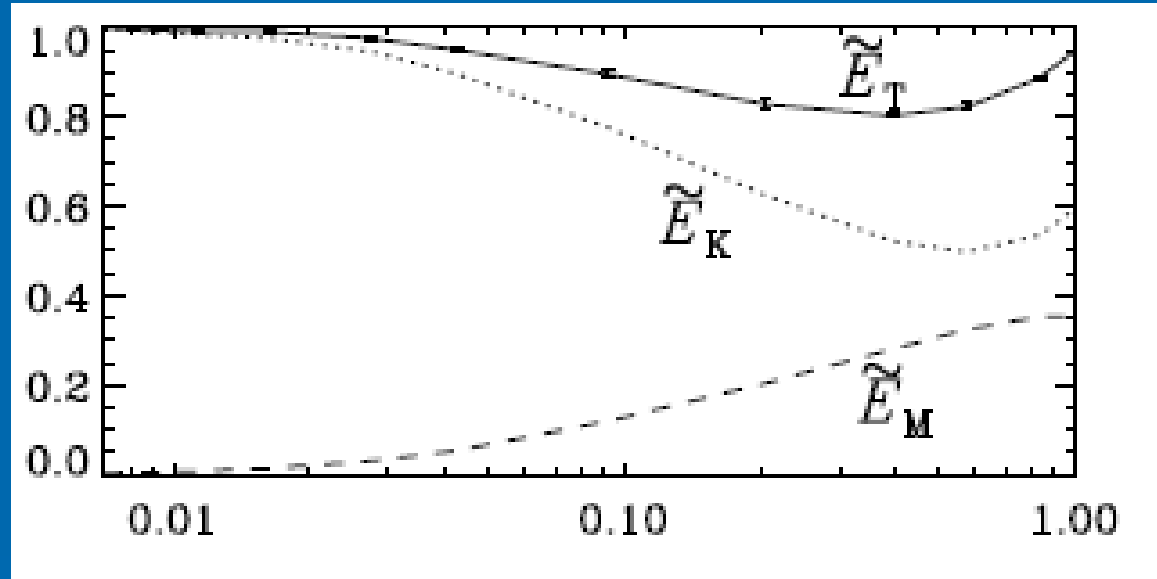
$$P_B = \frac{\bar{B}^2}{8\pi};$$

$$P_B^{eff} = \frac{\bar{B}^2}{8\pi} \left(1 - \frac{a + 2a_1}{3} \right)$$

$$\langle b^2 \rangle = a\bar{B}^2;$$

$$E_T^{(tot)} = E_T^{(tot)}(0) - \frac{a_1}{8\pi} \bar{B}^2;$$

$$P_{MHD} = \frac{\bar{B}^2}{8\pi} \left(1 - \frac{a + 2a_1}{3} \right) + \frac{2}{3} E_T^{(tot)}(0)$$

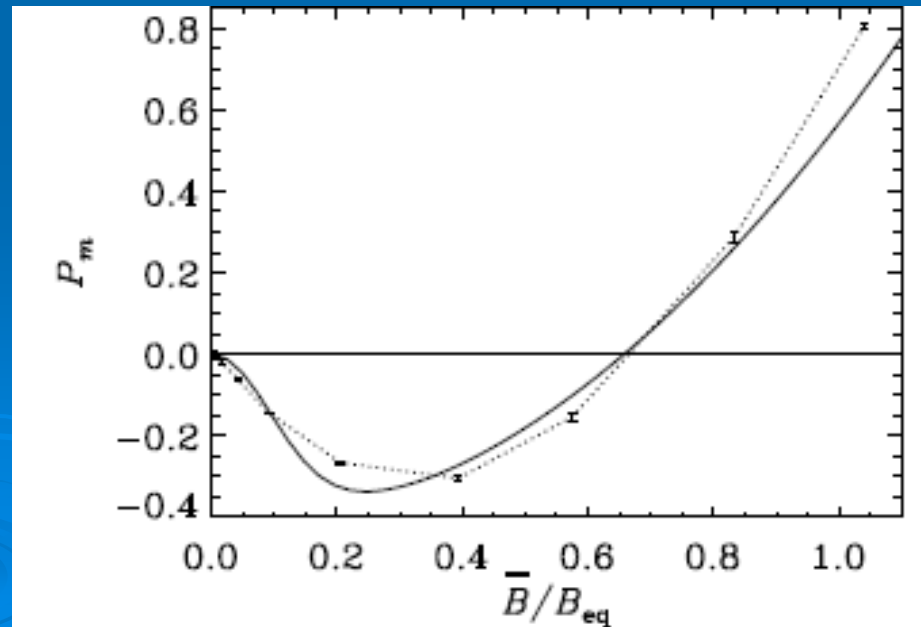
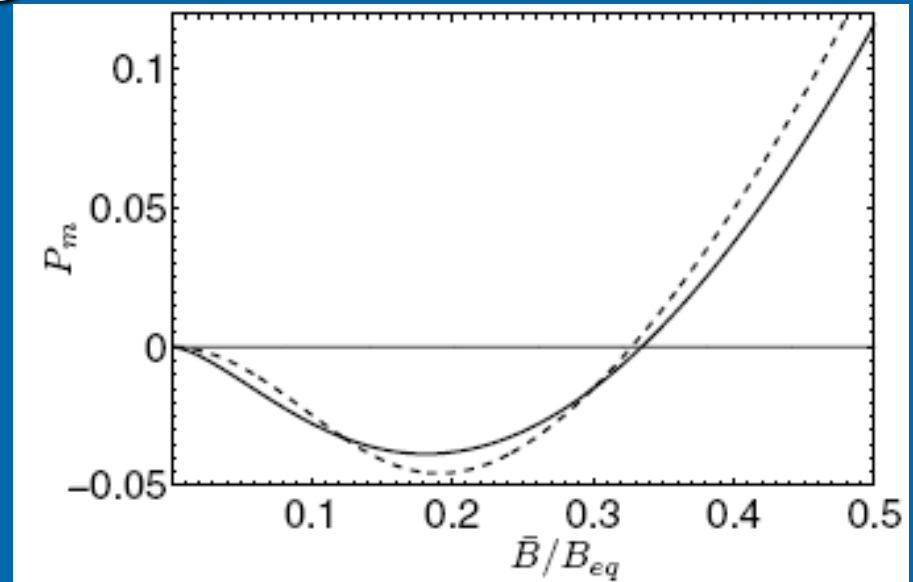


(Brandenburg, Kleeorin, Rogachevskii 2010)

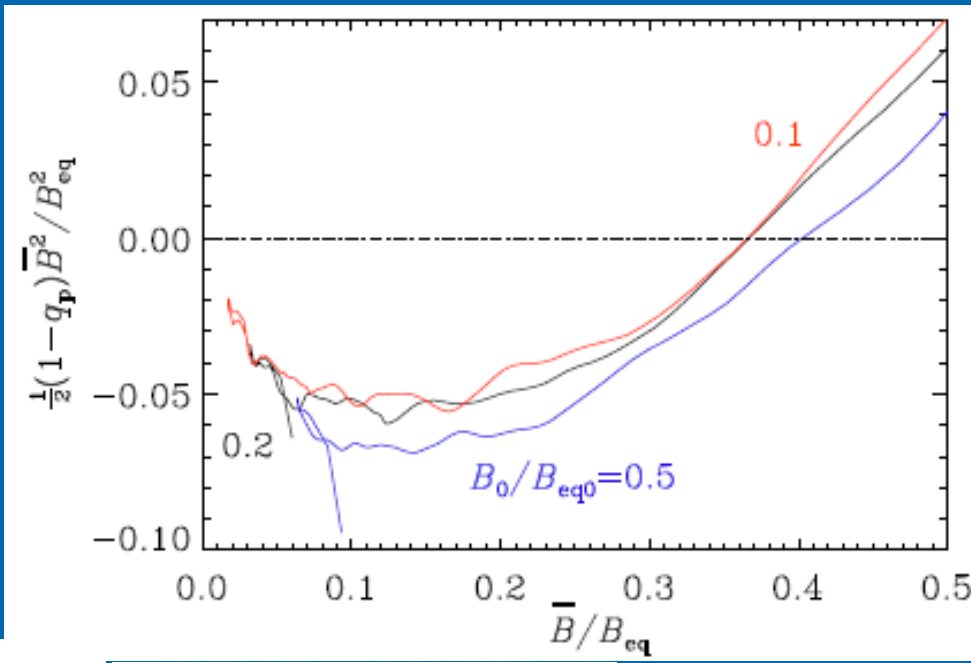
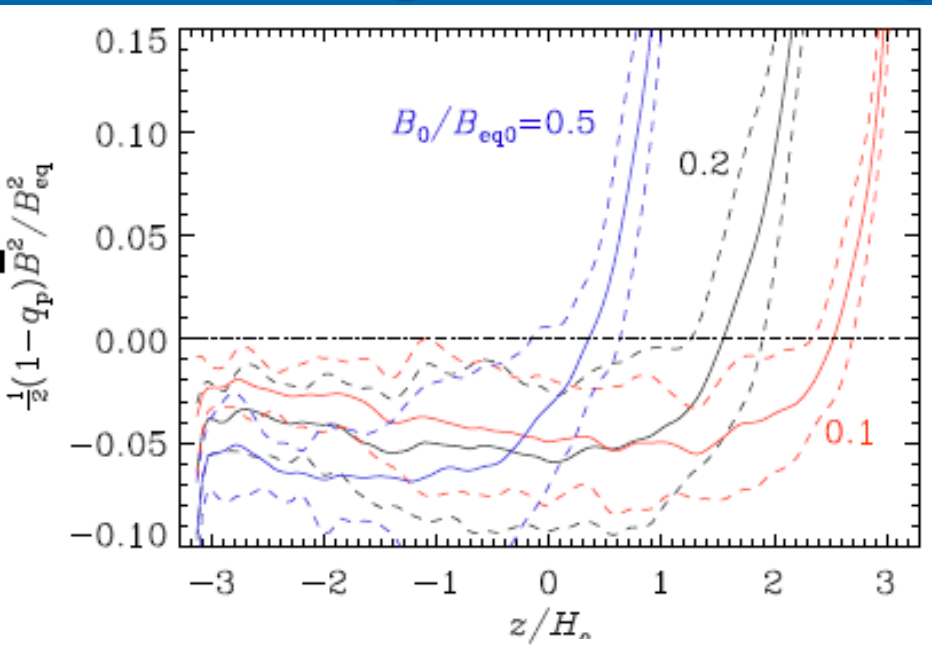
Negative Magnetic Pressure

$$P_m = (1 - q_p) \frac{\bar{B}^2}{B_{eq}^2}$$

(Brandenburg, Kleeorin,
Rogachevskii 2010)

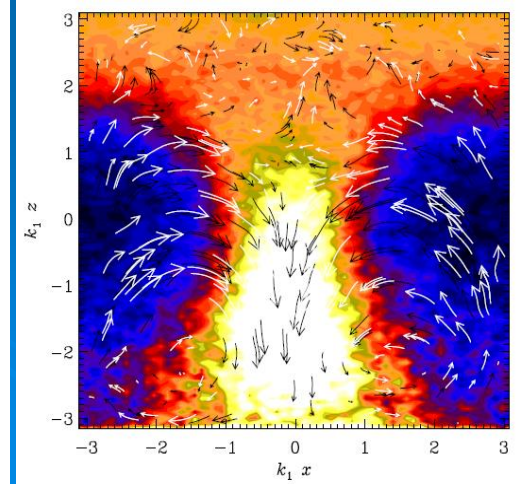
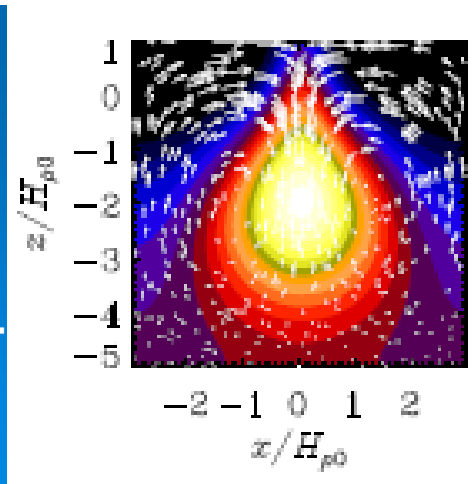


Negative Magnetic Pressure

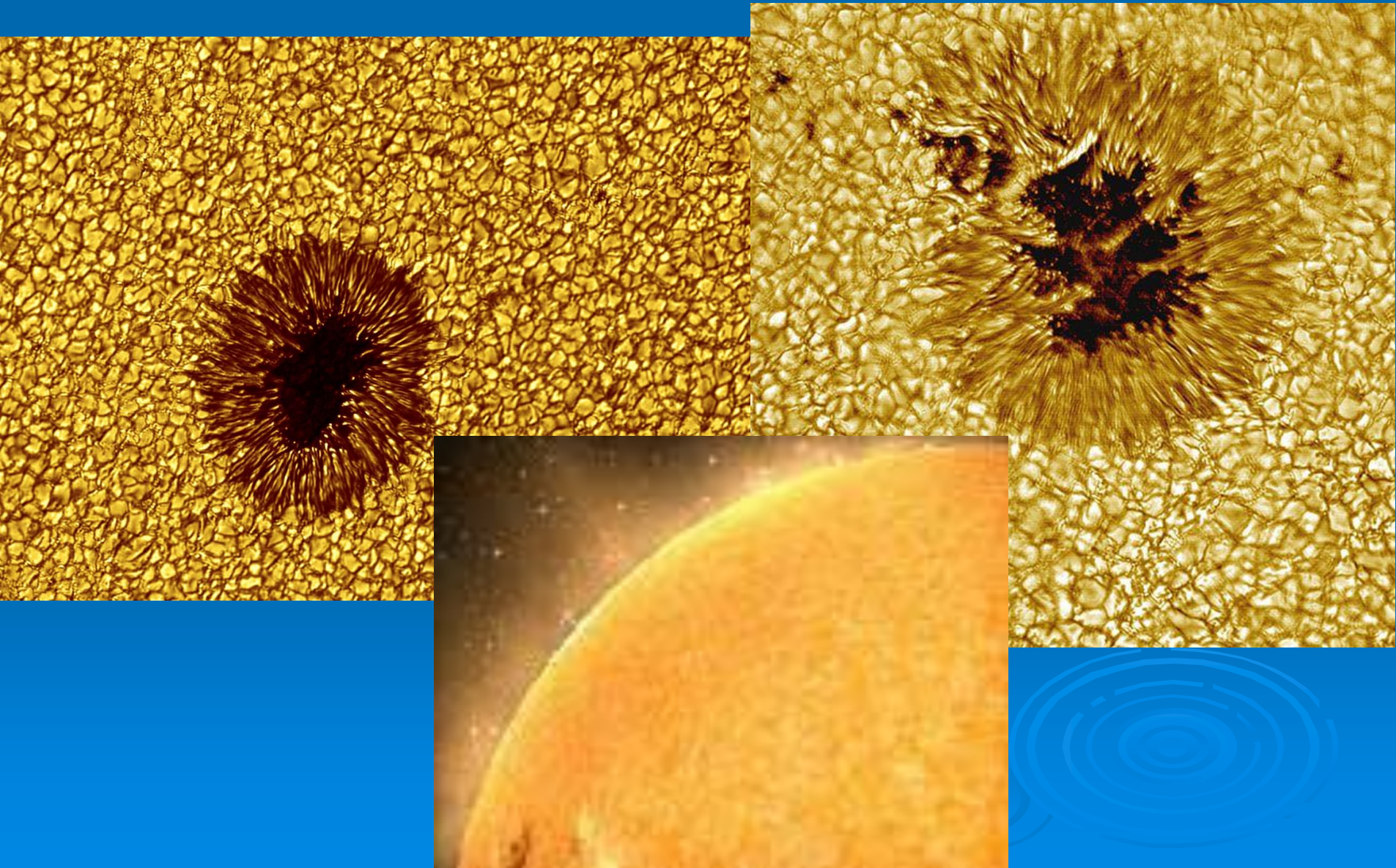


$$P_m = (1 - q_p) \frac{\overline{B^2}}{B_{eq}^2}$$

(Brandenburg & ell 2010-17)



Solar convection-examples



Nonlinear $\alpha\Omega$ dynamo

Mean magnetic field:

$$\mathbf{B} = \mathbf{e}_\varphi B(r, \vartheta) + \text{rot} \left[\mathbf{e}_\varphi A(r, \vartheta) \right]$$

$$r \sin(\vartheta) [\nabla \Omega \times \nabla A]_\varphi = \frac{1}{r} \left(\frac{\partial \Omega}{\partial r} \frac{\partial}{\partial \vartheta} - \frac{\partial \Omega}{\partial \vartheta} \frac{\partial}{\partial r} \right) (rA \sin(\vartheta))$$

$$\frac{\partial B}{\partial t} = rD \sin(\vartheta) [\nabla \Omega \times \nabla A]_\varphi + \Delta_s B$$

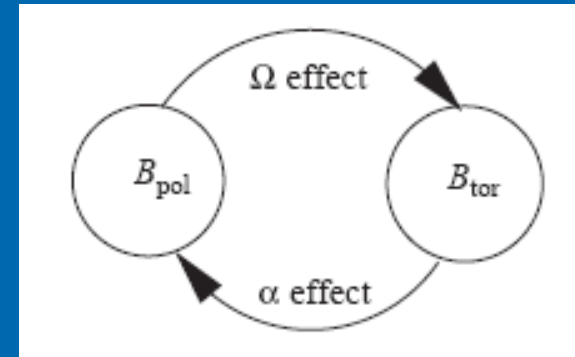
$$\frac{\partial A}{\partial t} = \alpha B + \Delta_s A; \quad \Delta_s = \Delta - (r \sin \vartheta)^{-2}$$

$$\frac{\partial \alpha_m}{\partial t} + \text{div} \mathbf{F} = \frac{Q}{4\pi\rho\eta_T} \left[-(\alpha_h + \alpha_m) \mathbf{B}^2 + \eta_T \mathbf{B} \cdot \text{rot} \mathbf{B} \right] - \frac{\alpha_m}{\tau_\chi}$$

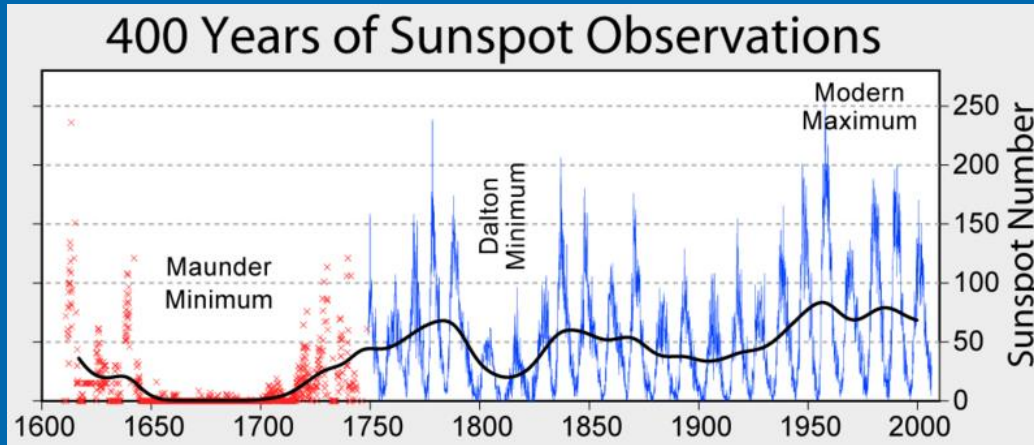
$$\tau_\chi \approx \tau_0 \text{Rm}$$

Kleeorin & Ruzmaikin (1982)

Reza Tavakol & Enrique Covas (1997)



RESULTS



$$\phi_m =$$

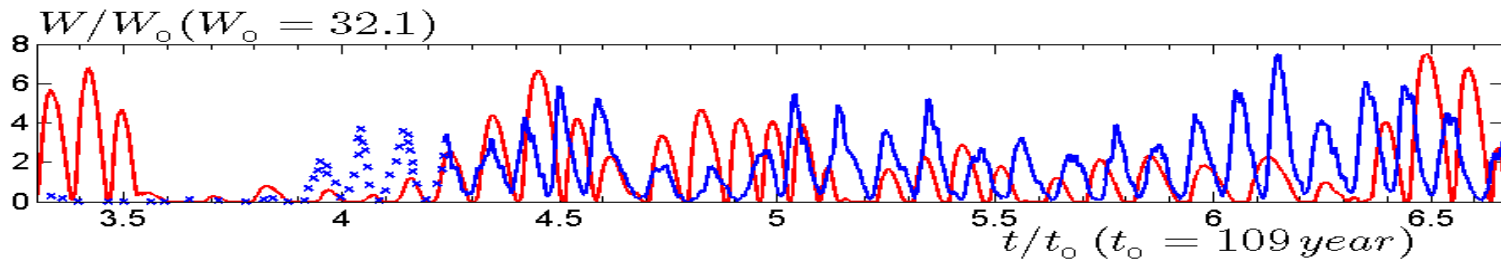
$$3 \left[1 - \arctan(\xi B_{tot}) / \xi B_{tot} \right] / \xi^2 B_{tot}^2 ;$$

$$D = -(4 \div 10) \times 10^3$$

$$\mu^2 \approx .1 \quad \kappa = .1$$

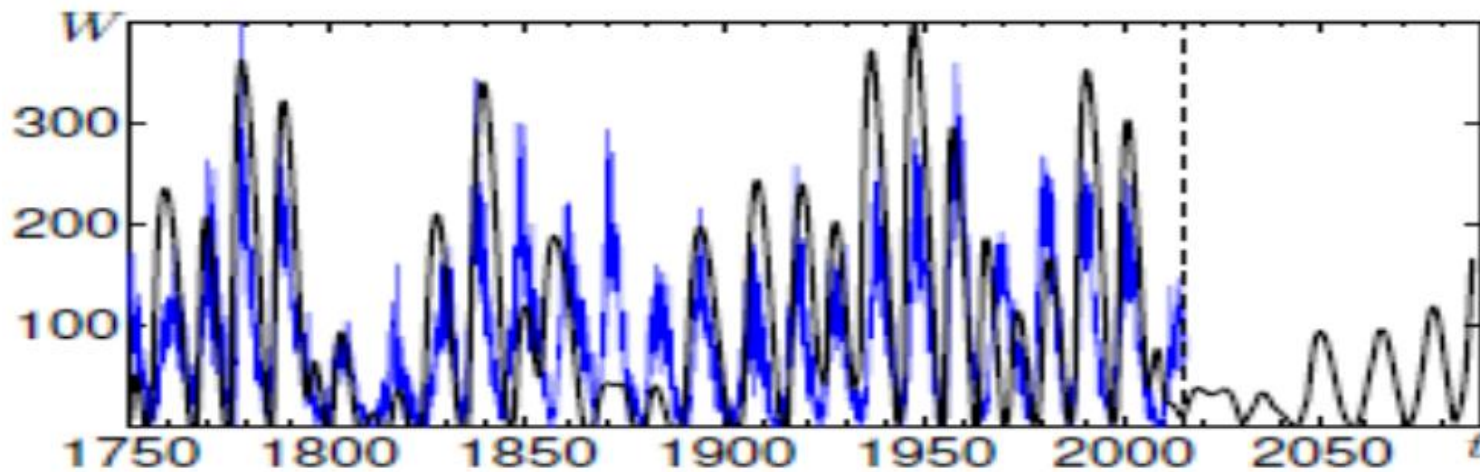
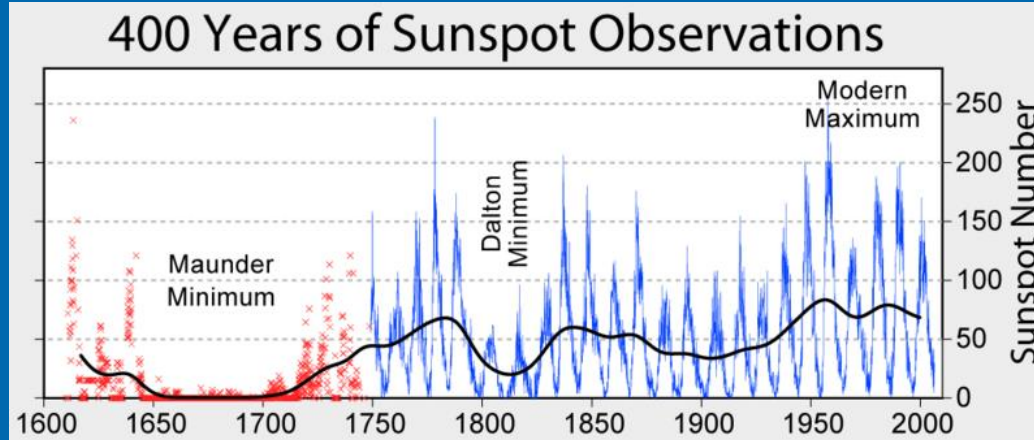
$$\sigma = 3 \quad T = 6.3$$

$$R_\alpha = 2 \quad \xi = l\sqrt{2}/R_\odot \approx .3$$

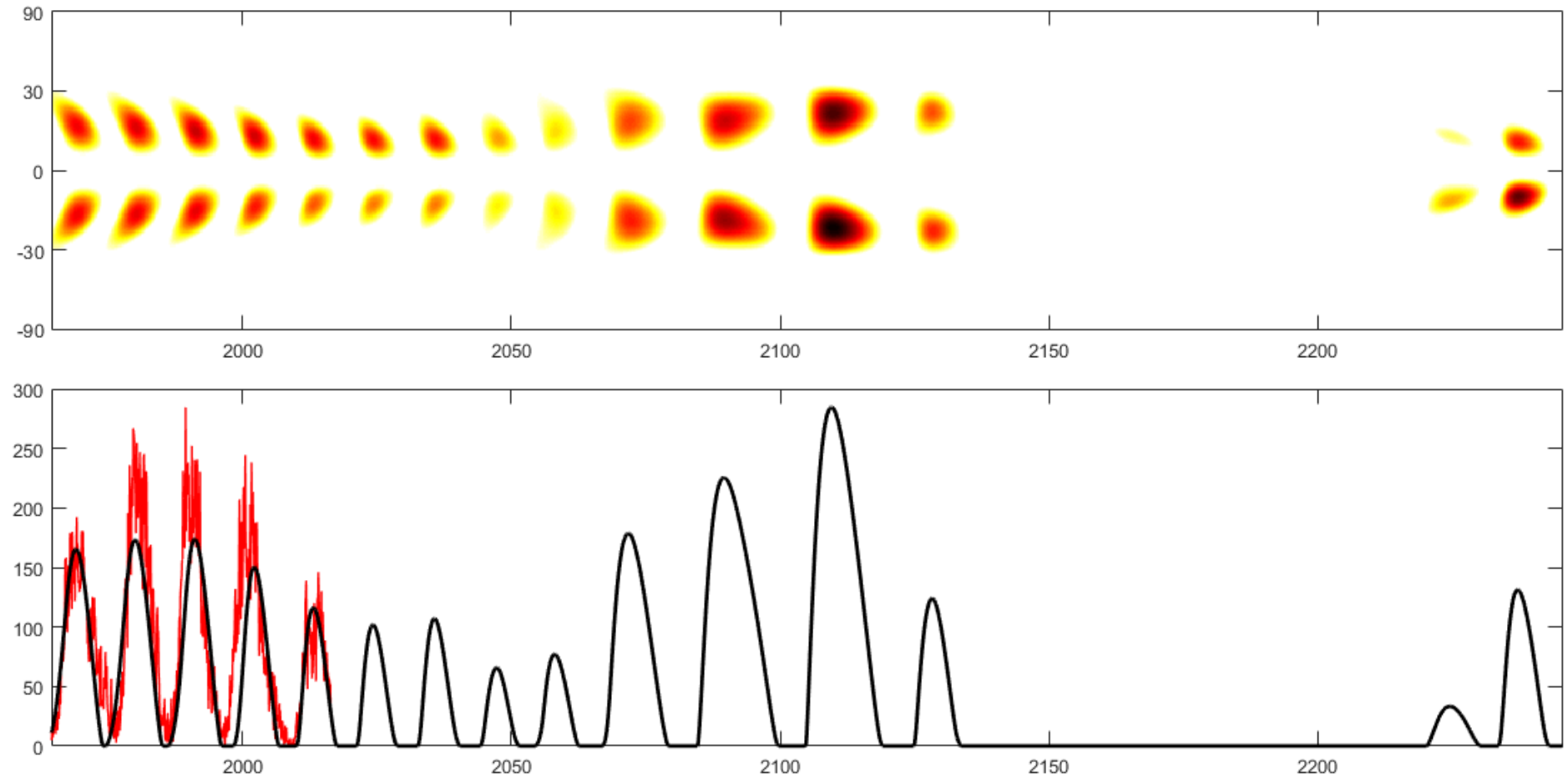


RESULTS

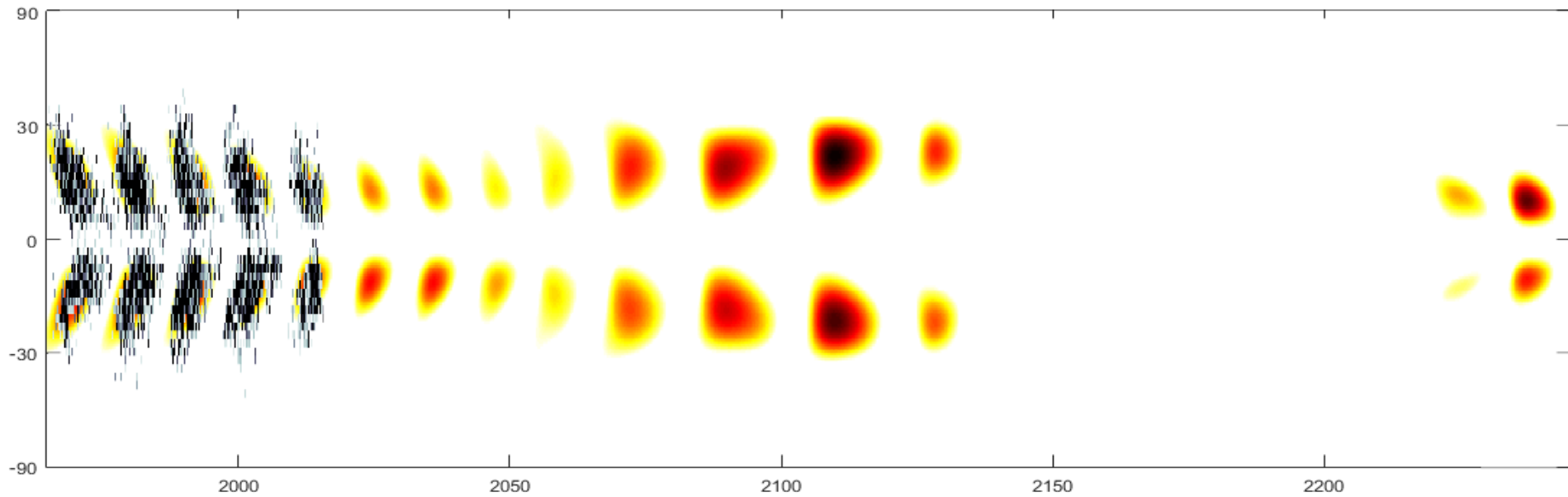
(Ya. Kleorin & ell. 2016)



RESULTS: 100-200 years prediction of solar activity



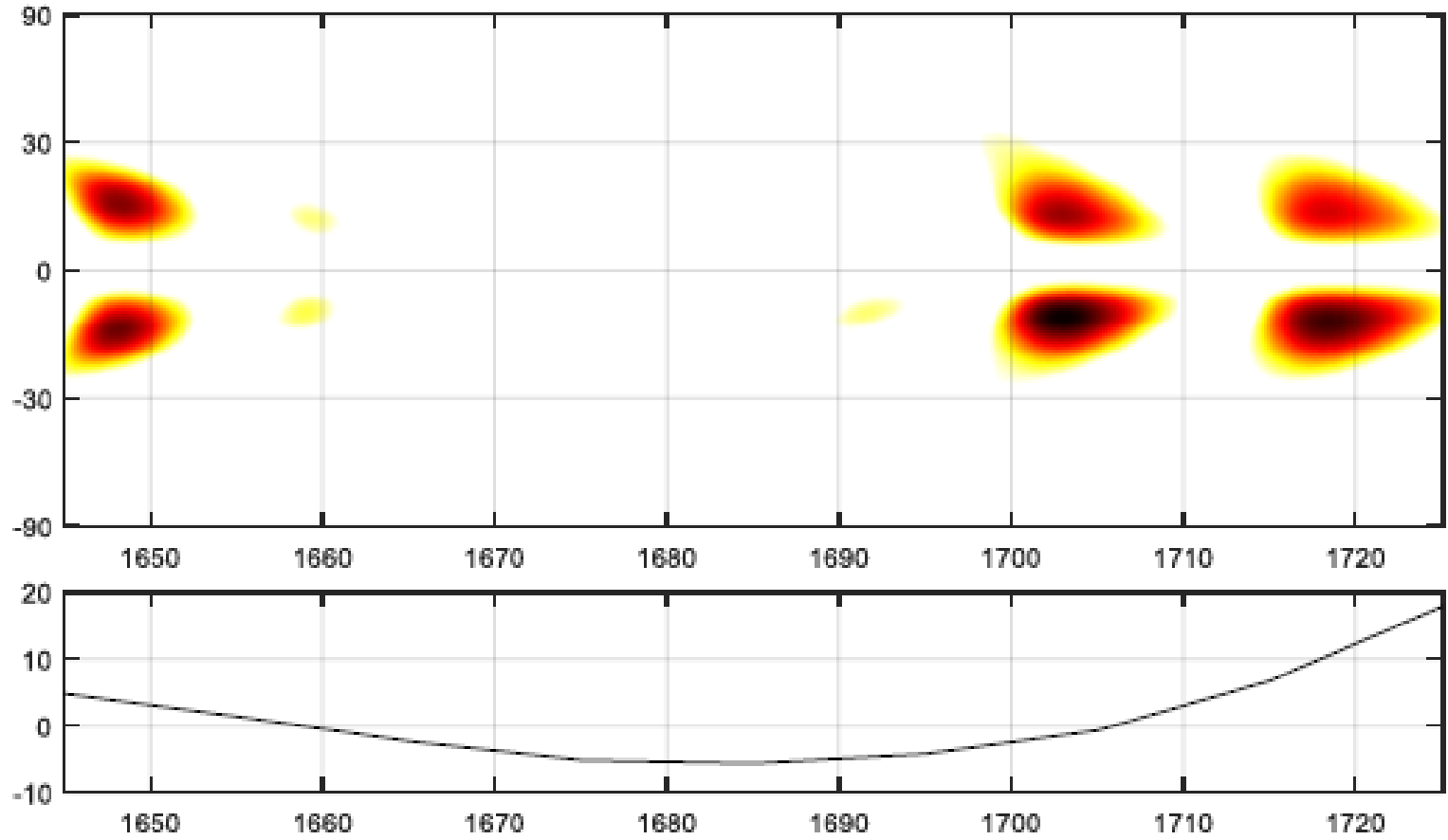
RESULTS: 100-200 years prediction of solar activity $R_{\odot}/L_{\odot} \approx \mu^{-1}$



(Safiullin & all 2018)

RESULTS: 100-200 years prediction of solar activity $R_{\odot}/L_{\odot} \approx \mu^{-1}$

Все изображения с 1645 г. до 1725 г. (минимум Маундера)



(Safiullin & all 2019)

STARS

$$B_s \approx B_{s2} \left(\sqrt{\frac{\bar{\Omega}_*^3}{\bar{\Omega}_{cr}^3}} - 1 \right)^{\frac{1}{2}}$$

Таблица 2

Спектрал. тип	F0	F5	G0	G2	G5	K0	K5	M0	M5
M_*/M_\odot	1.7	1.3	1.1	1	0.91	0.78	0.7	0.47	0.21
R_*/R_\odot	1.3	1.2	1.05	1	0.933	0.85	0.74	0.63	0.32
$\eta_T^{(*)}/\eta_T^{(\odot)}$	-----	1.89	1.31	1	0.74	0.48	0.45	0.2	$\approx .5 \div 4$
ρ_*/ρ_\odot	0.77	0.75	0.95	1	0.929	1.27	1.73	1.88	6.4
B_0 (Гаусс)	-----	75	67	55	51	41	35	23	217
\bar{T}	1.25	1.13	1.024	1	0.971	0.907	0.763	0.678	0.540
B_{s2} (Гаусс)	-----	23	56	74	118	180	345	3,560	$188 \times 10^3 - ?$
μ_*	∞	6.18	4	3.34	3.16	3.12	3.47	2.26	1
\bar{H}_*	0	0.649	0.877	1	0.986	0.909	0.712	0.695	1.069
$\bar{\Omega}_\sigma$	∞	1.238	0.470	0.276	0.210	0.160	0.245	0.114	0.543

STARS

$$B_s \approx B_{s2} \left(\sqrt{\frac{\bar{\Omega}_*^3}{\bar{\Omega}_{cr}^3}} - 1 \right)^{\frac{1}{2}}$$

Таблица 3

Звезда	Спектральный тип	$\bar{\Omega}_*$	B_s (кГс, наблюдения)	B_s (кГс, теория)
Солнце	G2 V	1	1.5	0.193
HD 115383	G0 V	5.2	1.0	0.335
HD 20630	G5 V	2.7	1.8	0.792
HD 131511	K1 V	2.82	1.7	1.5
HD 26965	K1 V	0.686	1.7	0.506
HD 185114	K1 V	0.933	1.36	0.651
G1 171.2 A	K5 Ve	13.7	2.8	7.0
EQ Vir	K5 Ve	6.51	2.5	4.0

THE END

